

The Bardeen-Petterson Effect in accreting supermassive black-hole binaries

Disc breaking and critical obliquity

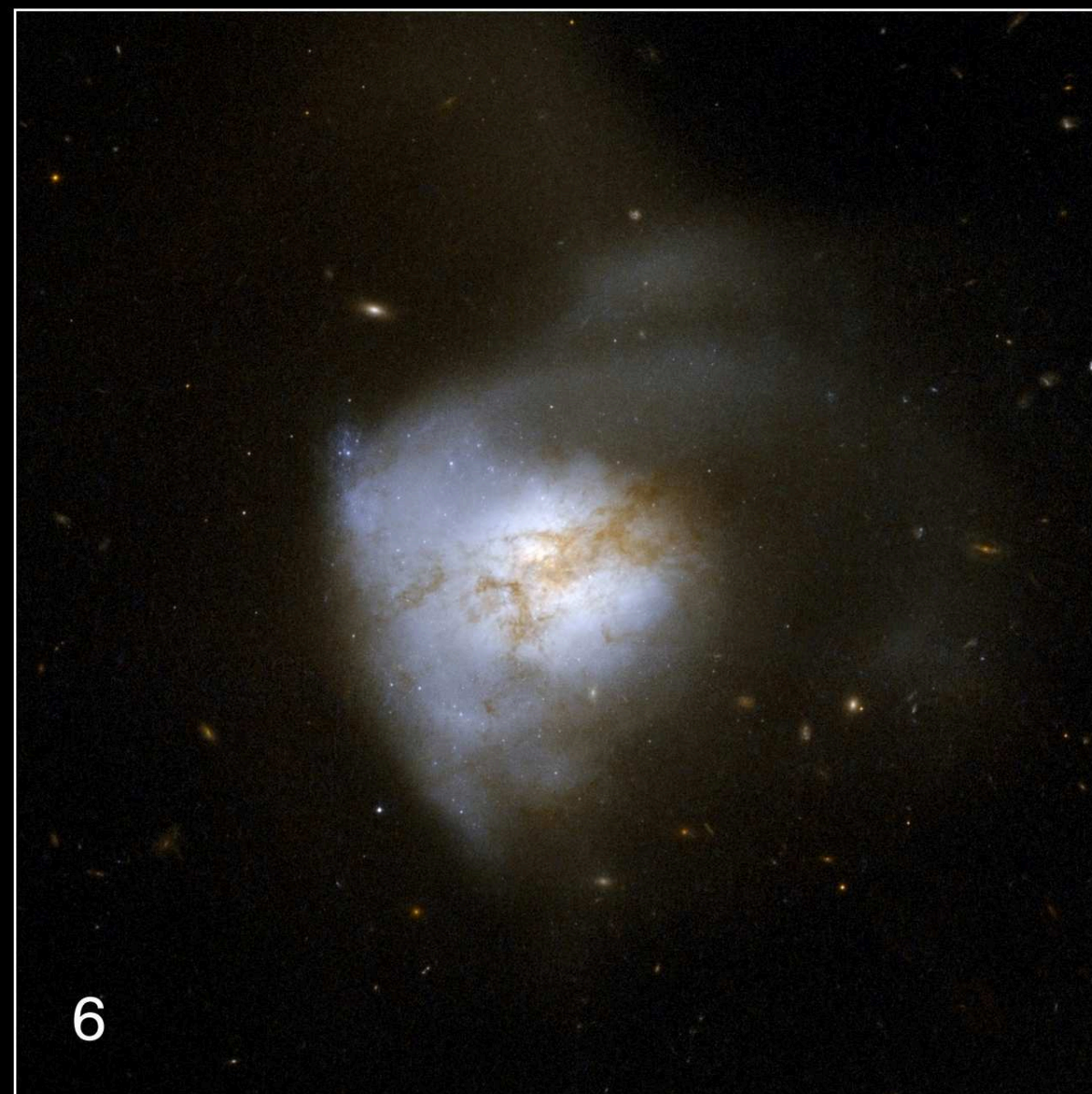
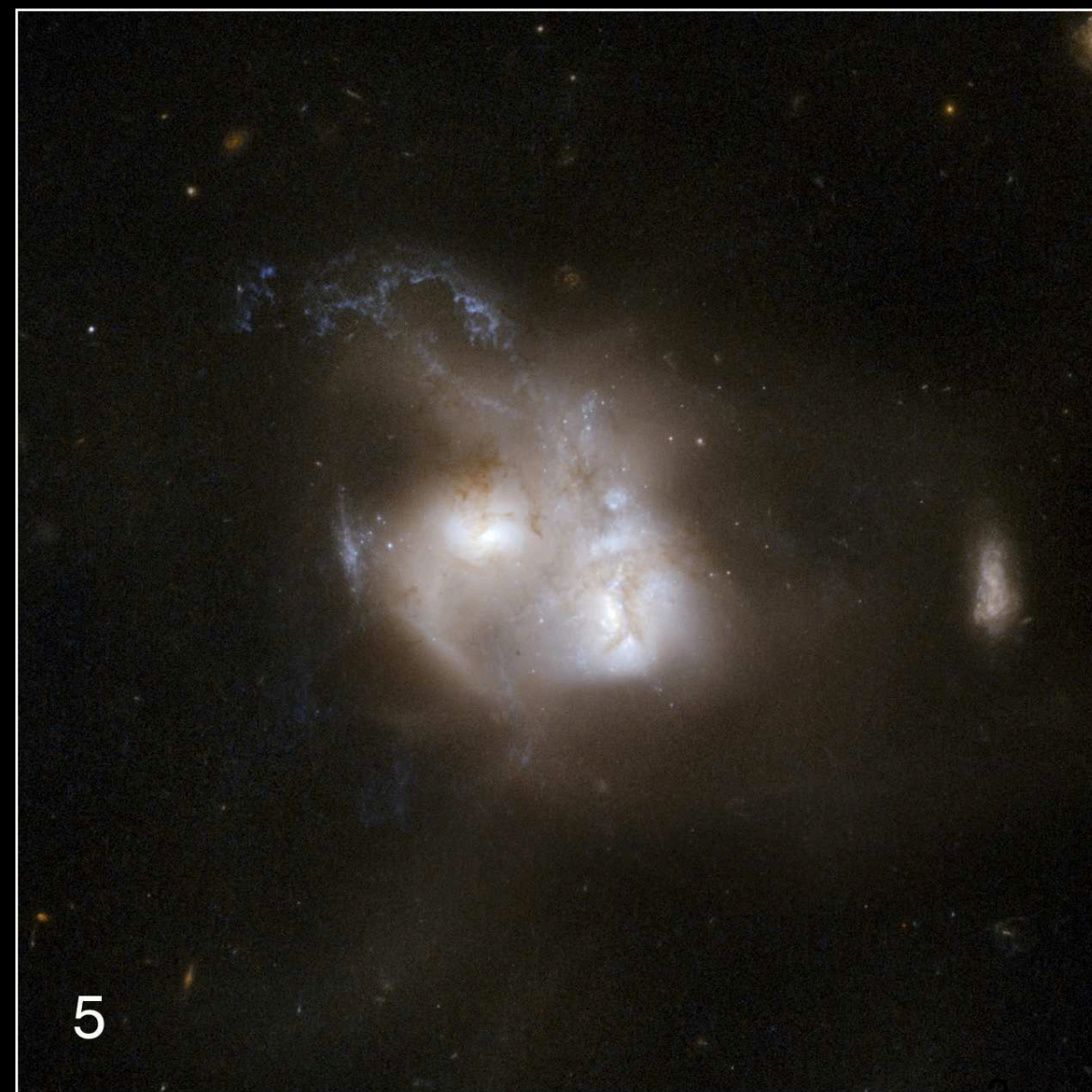
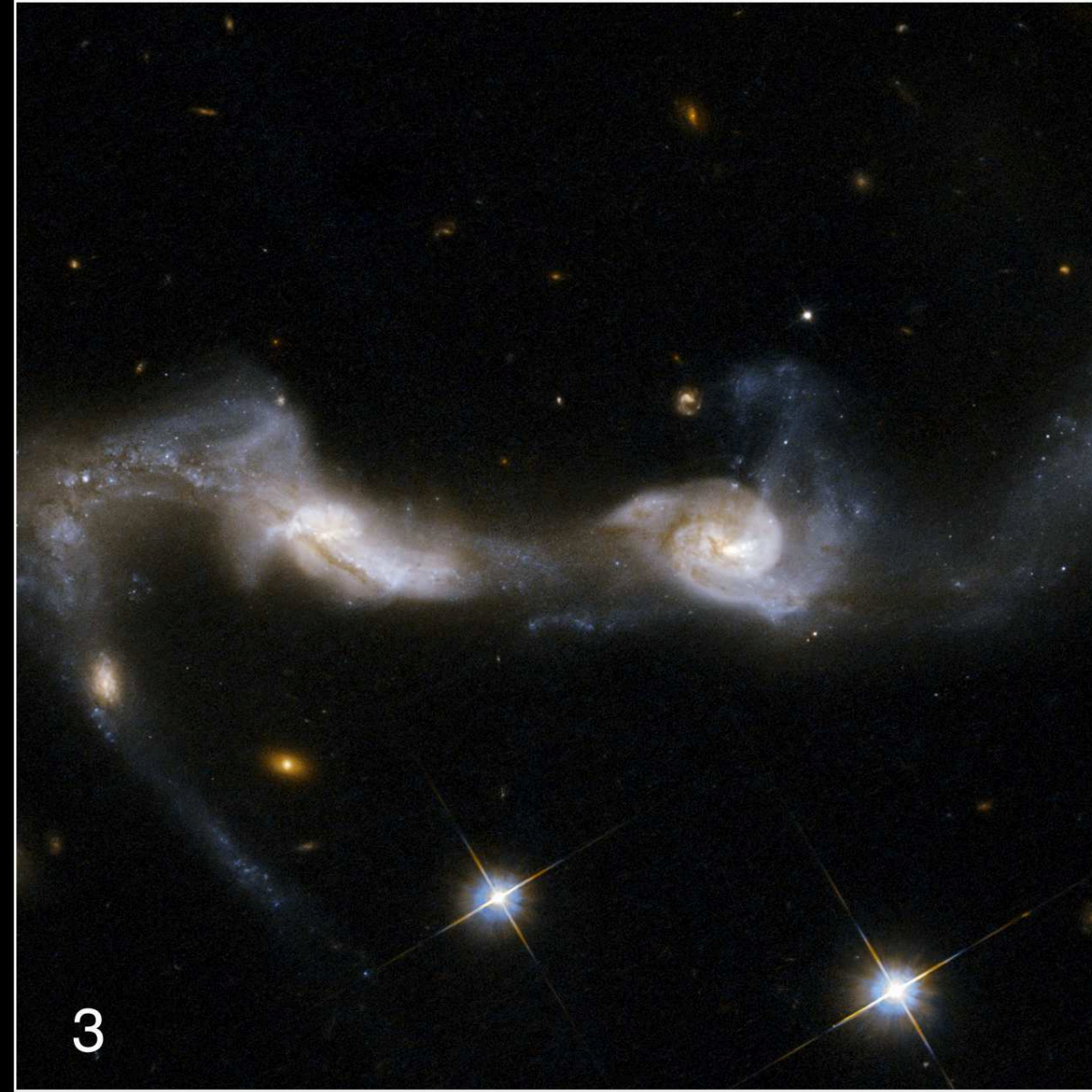
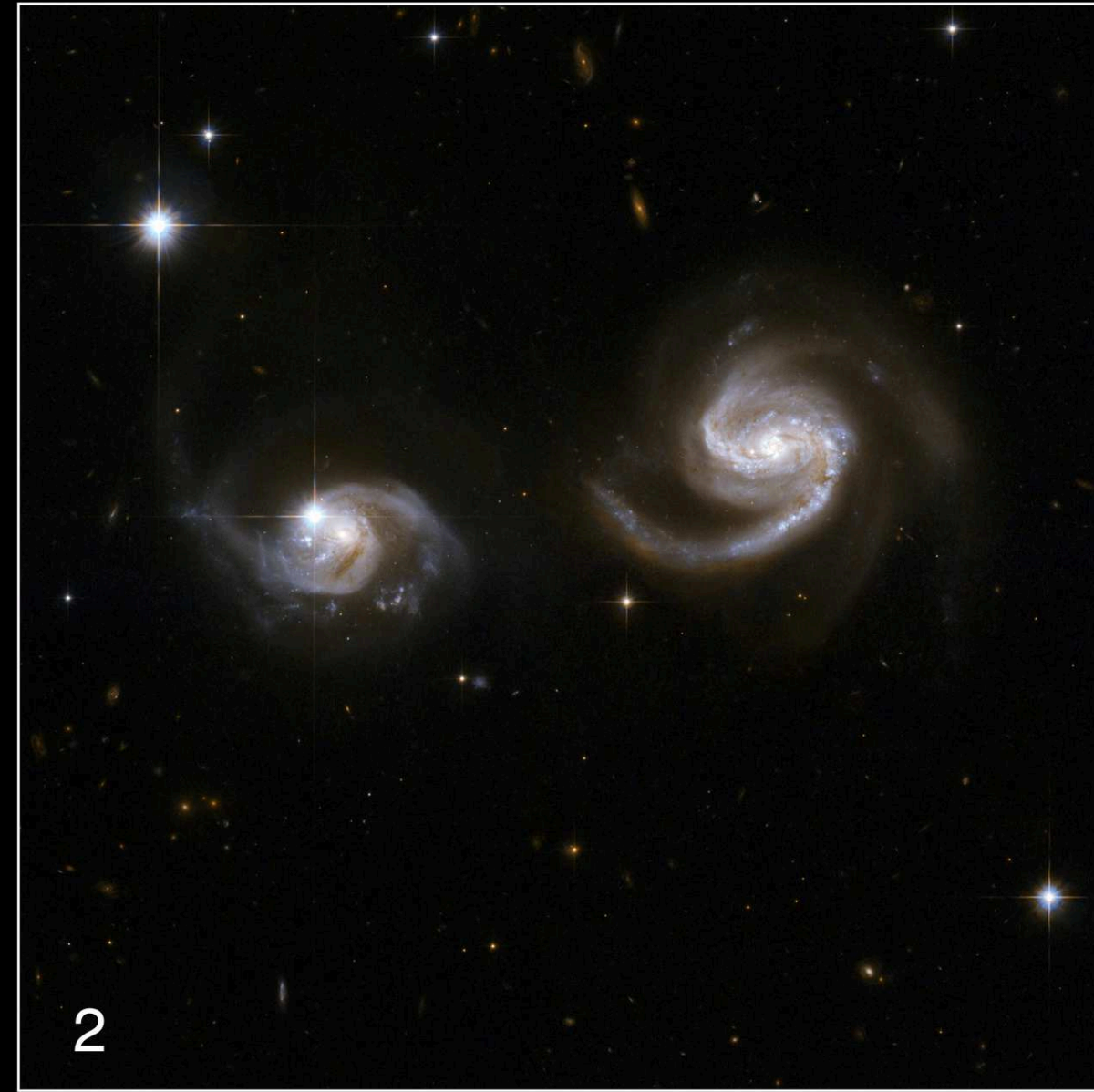
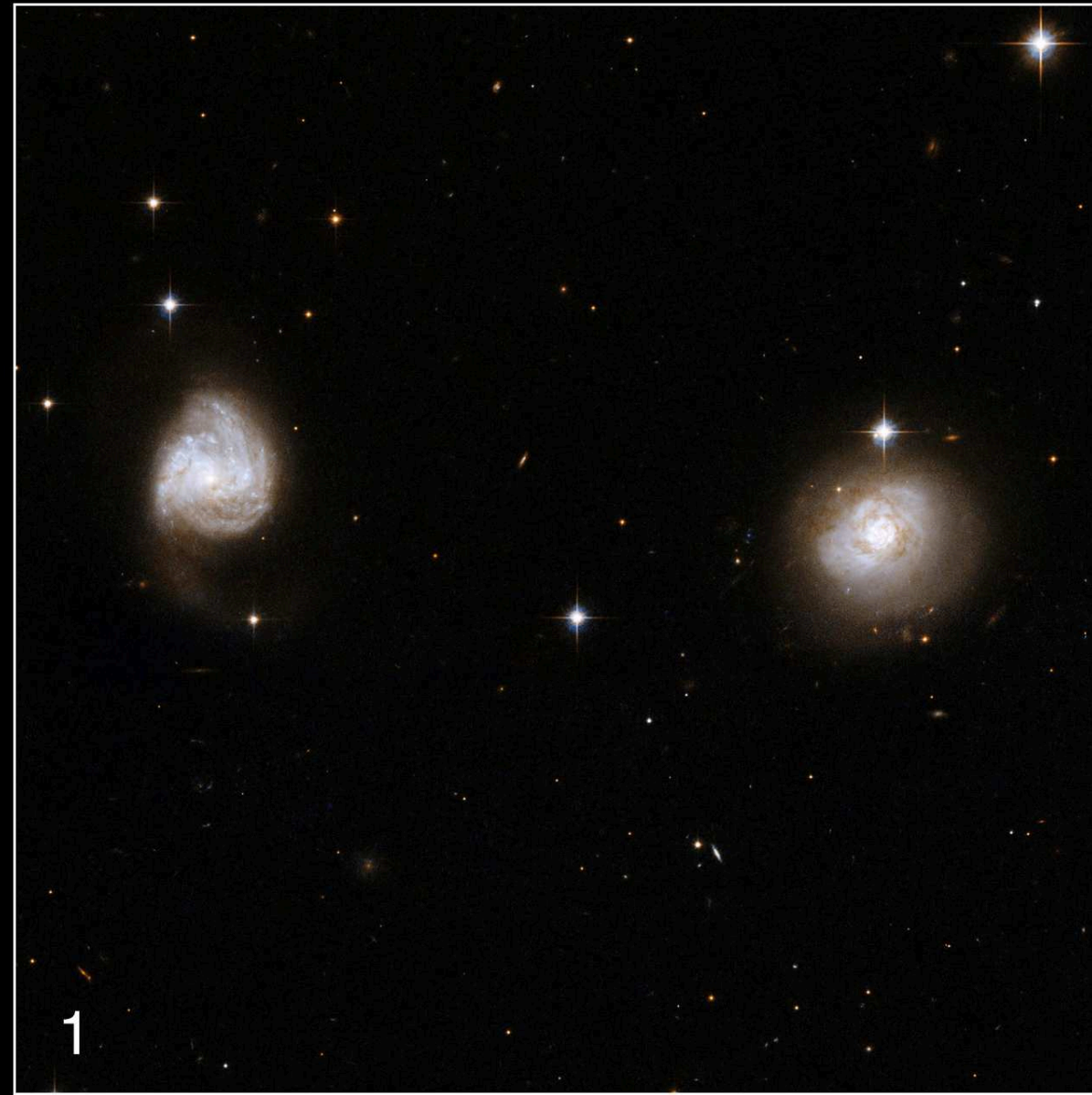


Rebecca Nealon, Stephen Hawking Fellow

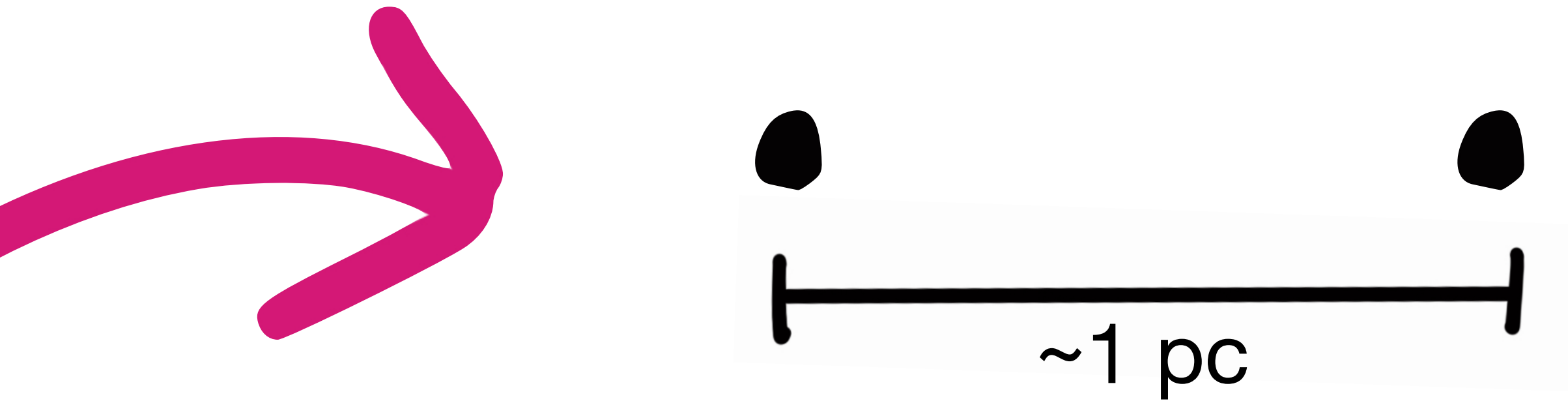
with Enrico Ragusa, Davide Gerosa, Giovanni Rosotti and Riccardo Barbieri



How do supermassive black holes merge?

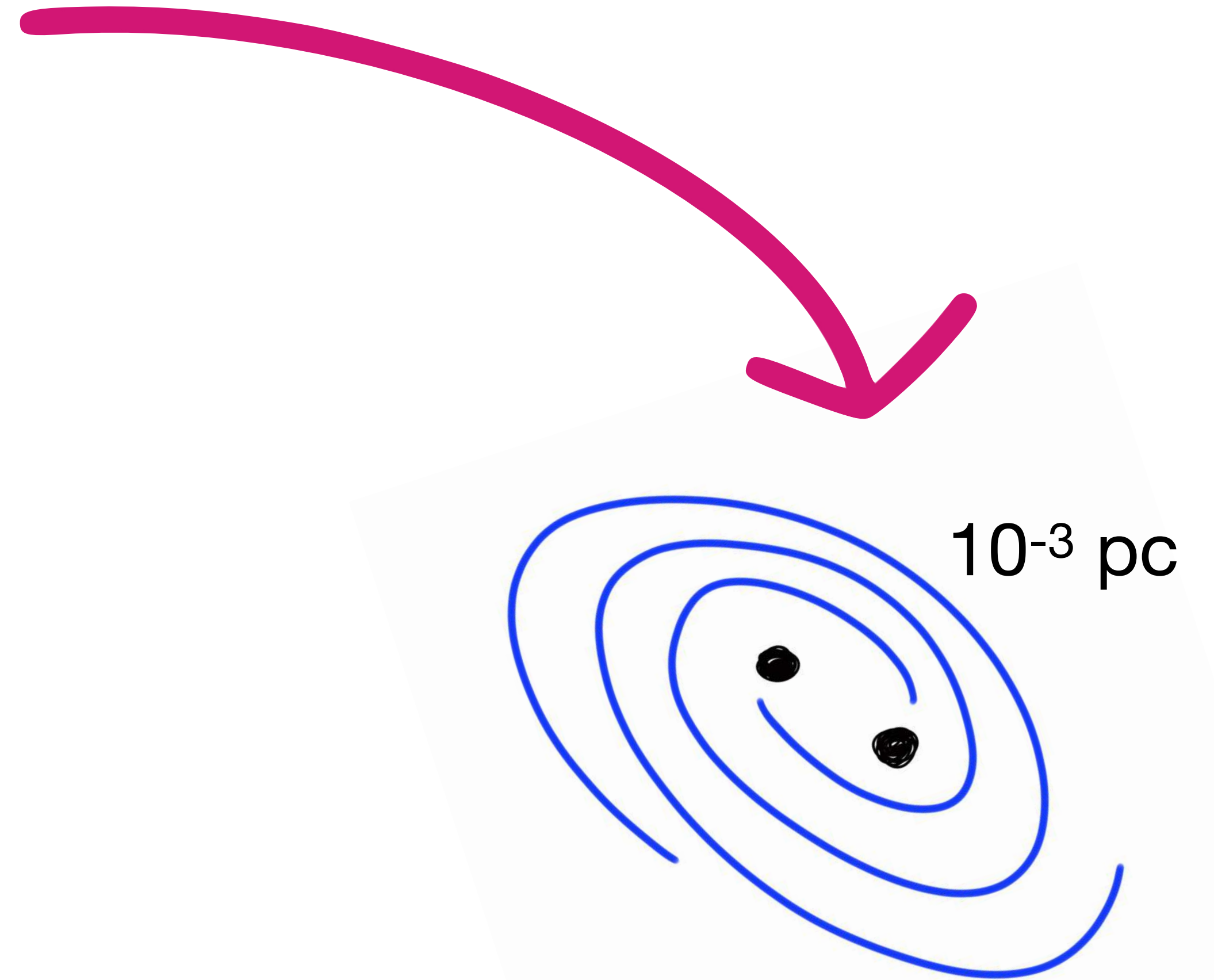


How do supermassive black holes merge?



The Final Parsec Problem

- Triaxial galactic potentials (Poon and Merritt 2004)
- Dynamical interactions in supermassive black hole triples (Bonetti et al. 2019)
- Gas accretion (e.g. Armitage and Natarajan 2002 and many others)

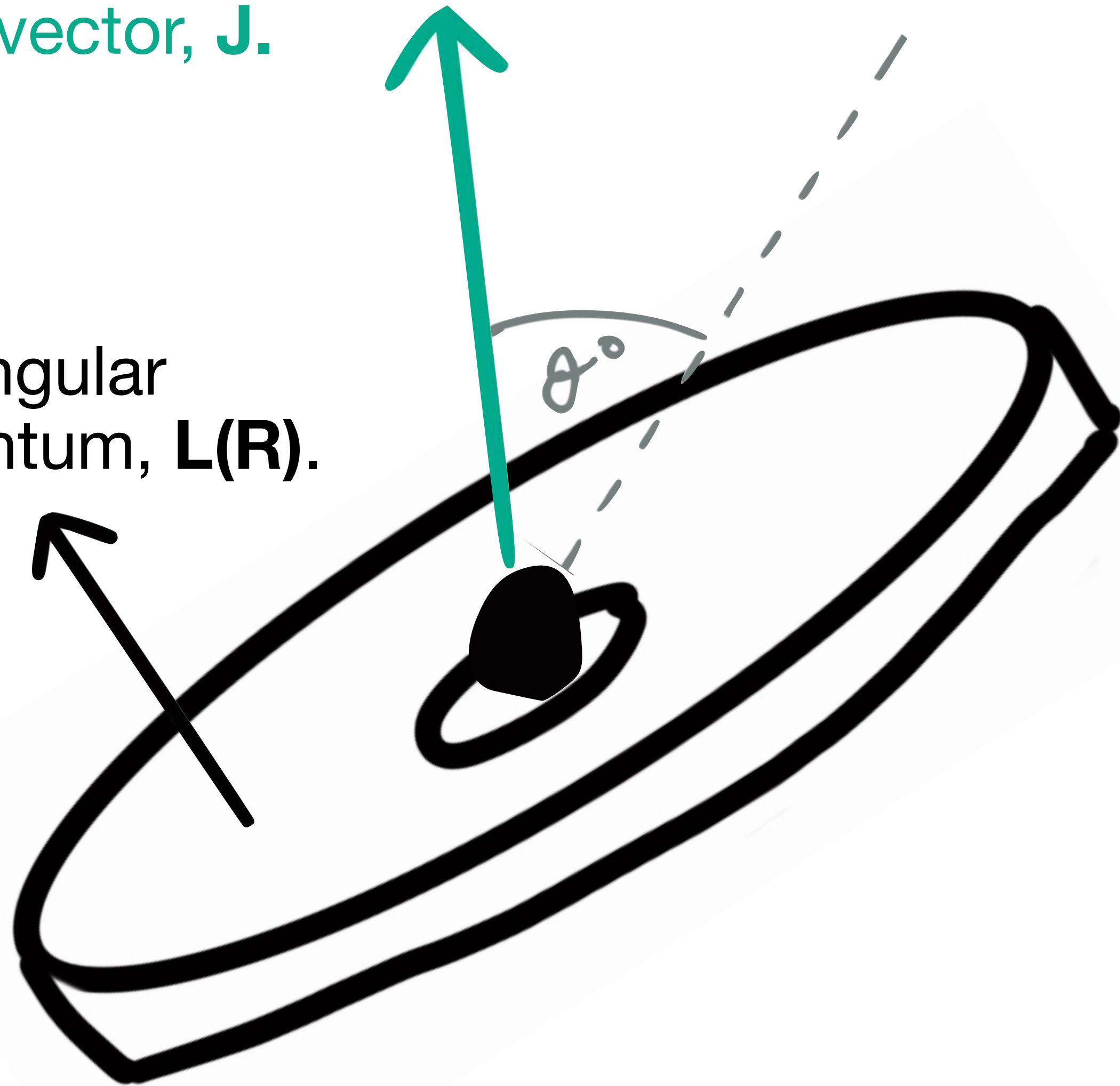


The emission of gravitational waves are responsible for the final merge

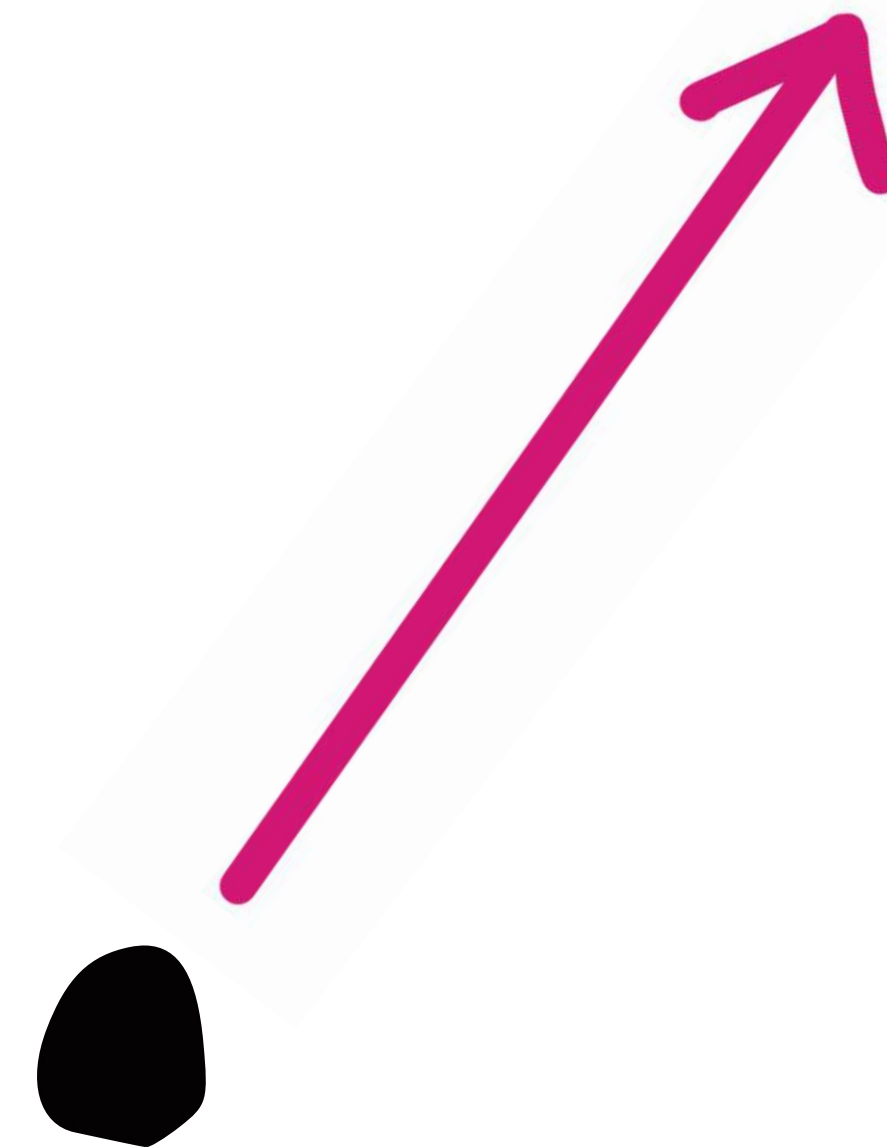
Disc driven migration

Primary black hole
spin vector, \mathbf{J} .

Disc angular
momentum, $\mathbf{L}(\mathbf{R})$.



Orbital angular momentum of
secondary, \mathbf{L}_* .

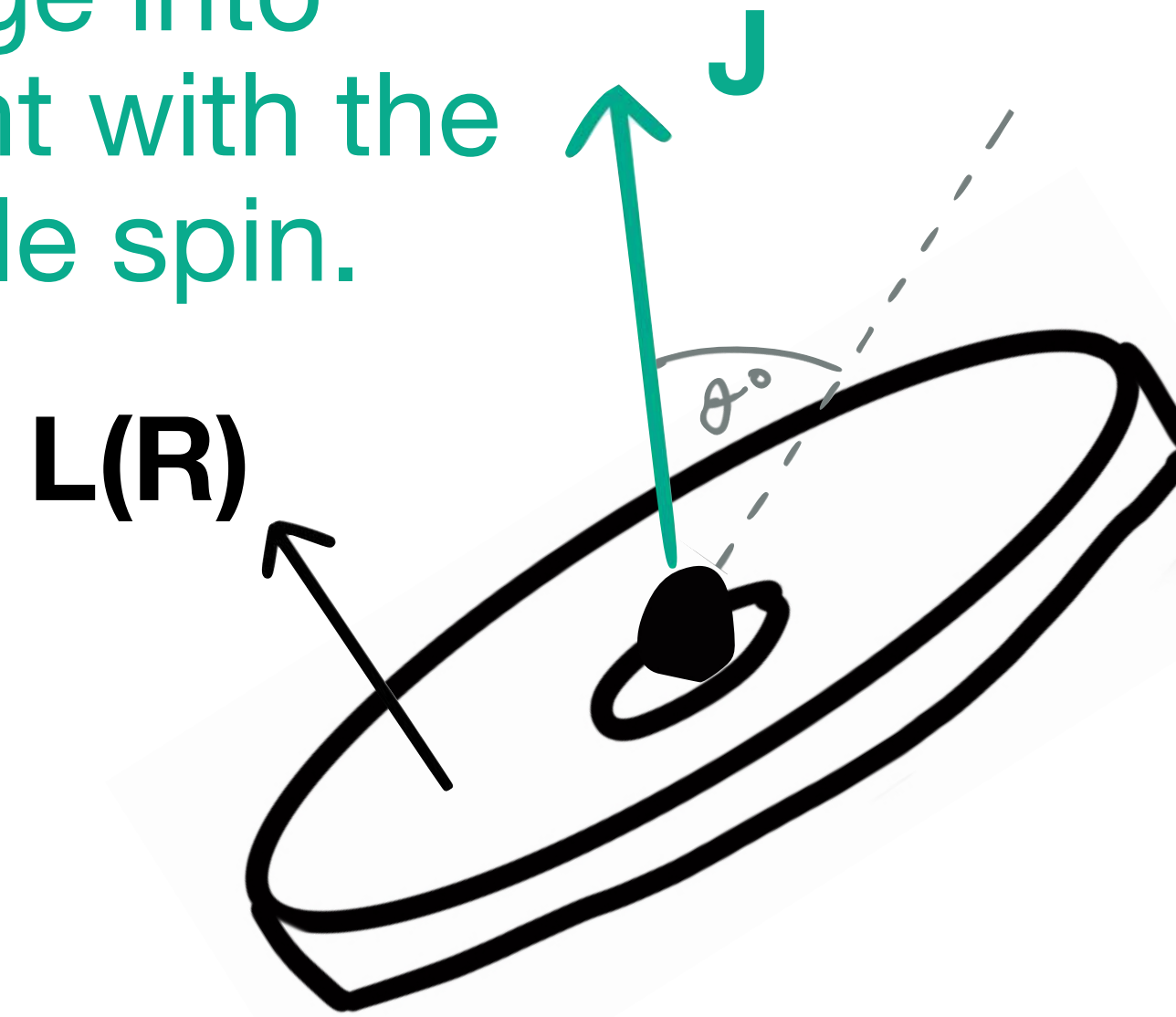


Disc driven migration

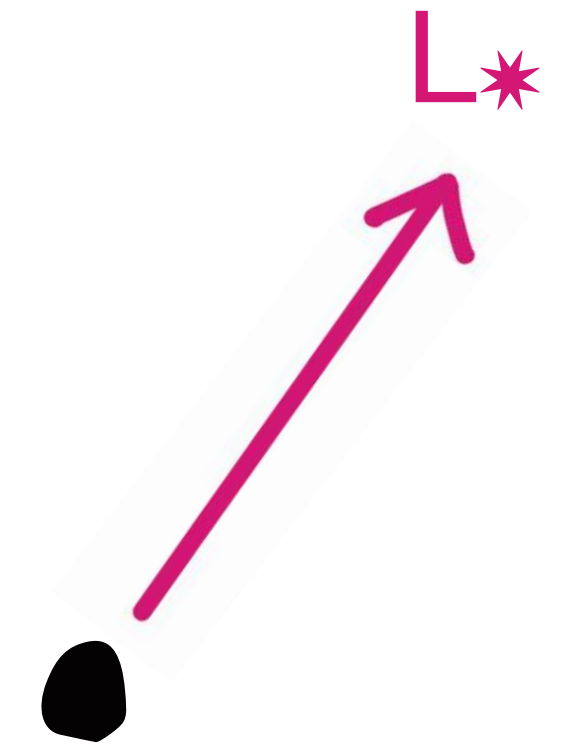
$$\frac{d \cos \theta}{dt} = \frac{d\hat{\mathbf{J}}}{dt} \cdot \hat{\mathbf{L}}_*$$

$$\frac{d\mathbf{J}}{dt} = - \int_{R_{\min}}^{R_{\max}} \frac{2G}{c^2} \frac{\mathbf{J} \times \mathbf{L}}{R^3} 2\pi R dR$$

Tries to pull the inner edge into alignment with the black hole spin.



Exerts a tidal torque on the outer edge of the disc.



1. What shape does the disc take?
2. How does this affect the alignment of the black holes?

A semi-analytic model

$$\frac{\partial \sigma}{\partial r} = - \left(\beta + \frac{1}{2} \right) \frac{\sigma}{r} - \frac{\zeta \sigma \psi^2}{3r} \frac{\tilde{\alpha}_2(\alpha, \psi)}{\tilde{\alpha}_1(\alpha, \psi)} + \frac{r^{-\beta-1}}{3\tilde{\alpha}_1(\alpha, \psi)} - \frac{\sigma}{\tilde{\alpha}_1(\alpha, \psi)} \frac{\partial \tilde{\alpha}_1(\alpha, \psi)}{\partial r},$$

$$\frac{\partial^2 \hat{\mathbf{L}}}{\partial r^2} = \frac{\partial \hat{\mathbf{L}}}{\partial r} \left[- \frac{2r^{-\beta-1}}{\zeta \tilde{\alpha}_2(\alpha, \psi) \sigma} + \frac{3 \tilde{\alpha}_1(\alpha, \psi)}{\zeta r \tilde{\alpha}_2(\alpha, \psi)} - \left(\beta + \frac{3}{2} \right) \frac{1}{r} - \frac{1}{\sigma} \frac{\partial \sigma}{\partial r} - \frac{1}{\tilde{\alpha}_2(\alpha, \psi)} \frac{\partial \tilde{\alpha}_2(\alpha, \psi)}{\partial r} \right] - \frac{\psi^2}{r^2} \hat{\mathbf{L}} - \left(\frac{R_{\text{LT}}}{R_0} \right) \frac{r^{-\beta-3}}{\tilde{\alpha}_2(\alpha, \psi)} (\hat{\mathbf{J}} \times \hat{\mathbf{L}}) - \left(\frac{R_{\text{tid}}}{R_0} \right)^{-7/2} \frac{r^{-\beta+3/2}}{\tilde{\alpha}_2(\alpha, \psi)} (\hat{\mathbf{L}} \cdot \hat{\mathbf{L}}_\star) (\hat{\mathbf{L}} \times \hat{\mathbf{L}}_\star).$$

Change of variables ...

$$r = \frac{R}{R_0}, \quad \sigma = \frac{2\pi}{\dot{M}} \alpha \nu_0 \Sigma;$$

and using

$$R_{\text{LT}} = \frac{4G^2 M^2 \chi}{c^3 \alpha \nu_0 \zeta},$$

$$R_{\text{tid}} = \left(\frac{2}{3} \frac{\sqrt{GM}}{GM_\star} R_\star^3 \alpha \nu_0 \zeta \right)^{2/7}$$

Aligns the inner disc with the black hole spin

Aligns the outer disc with the binary orbit

A semi-analytic model

Where the Lense-Thirring torque most strongly affects the warp profile

$$R_{\text{LT}} = \frac{4G^2 M^2 \chi}{c^3 \alpha v_0 \zeta},$$

Where the tidal external torque most strongly affects the warp profile

$$R_{\text{tid}} = \left(\frac{2}{3} \frac{\sqrt{GM}}{GM_\star} R_\star^3 \alpha v_0 \zeta \right)^{2/7}$$

Also introduce the convenient parameter:

$$\kappa = \left(\frac{R_{\text{tid}}}{R_{\text{LT}}} \right)^{-7/2}$$

The non-dimensional parameter κ

$$\kappa \simeq 0.66 \left(\frac{M}{10^7 M_{\odot}} \right)^2 \left(\frac{\chi}{0.5} \right)^2 \left(\frac{M_{\star}}{10^7 M_{\odot}} \right) \left(\frac{R_{\star}}{0.1 \text{pc}} \right)^{-3} \\ \times \left(\frac{H_0/R_0}{0.002} \right)^{-6} \left(\frac{\alpha}{0.2} \right)^{-3} \left[\frac{\zeta}{1/(2 \times 0.2^2)} \right]^{-3},$$

Measures the relative importance of the binary.

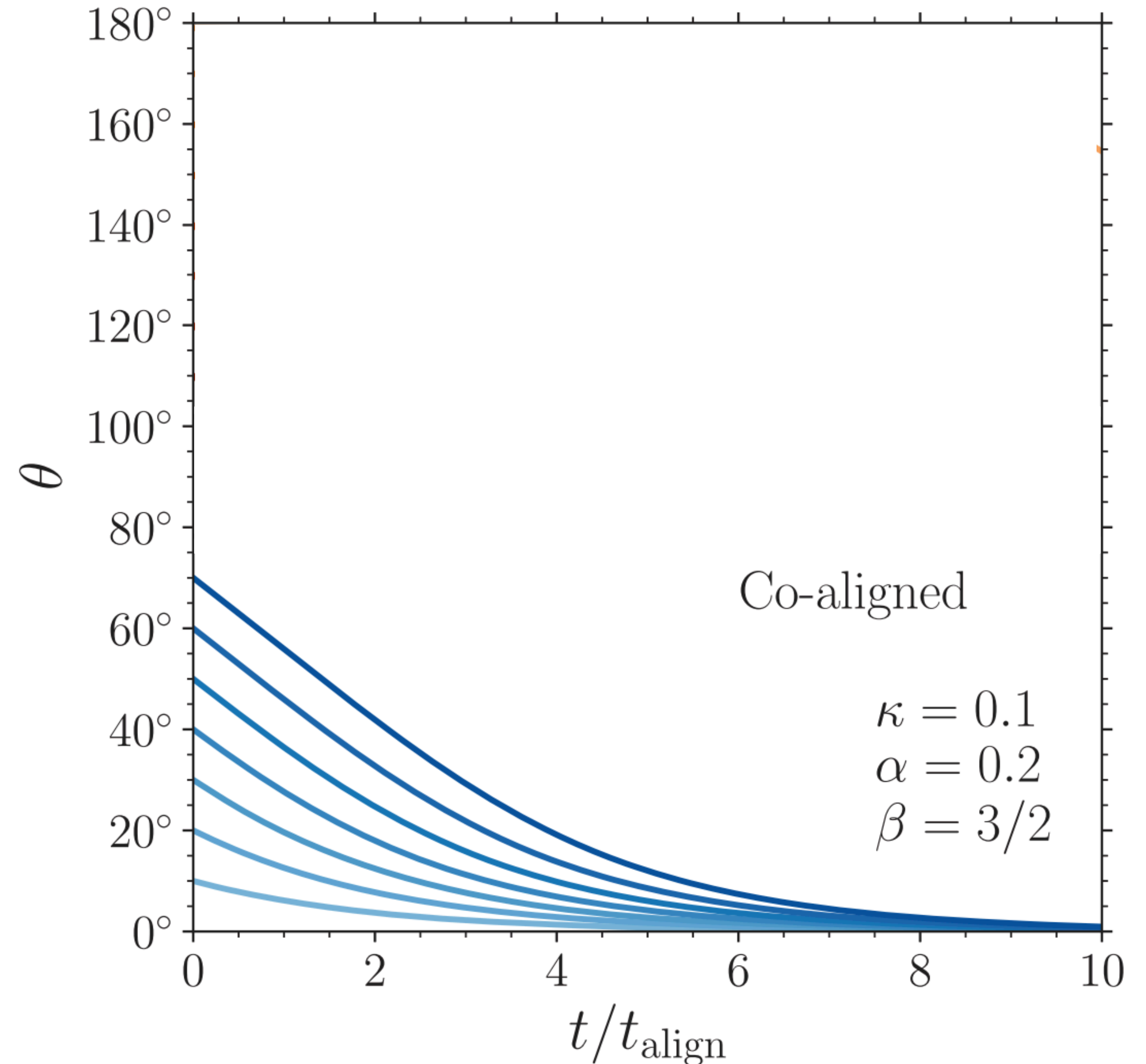
Large κ means the binary is very important.

Small κ means the binary is not important at all.

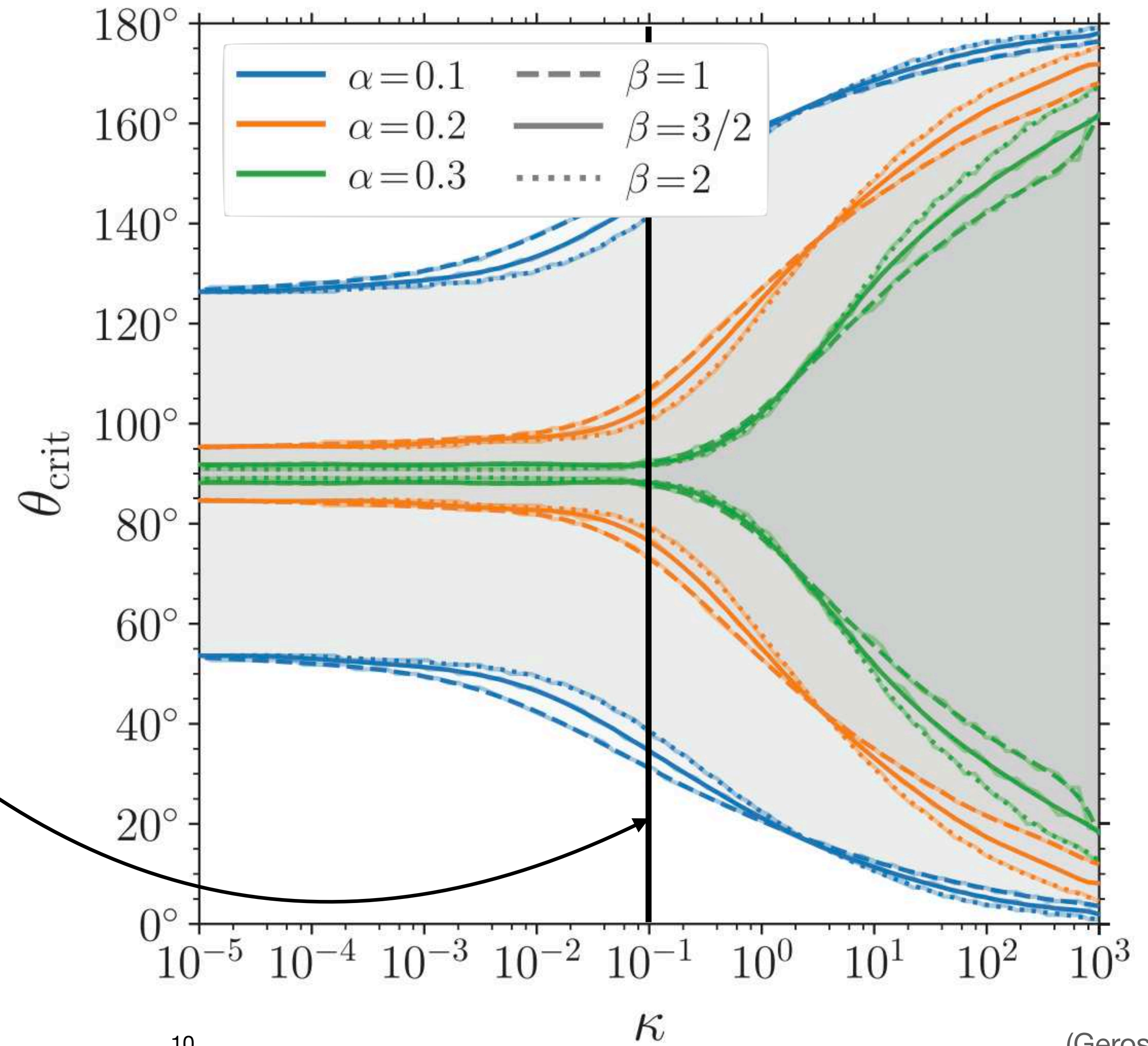
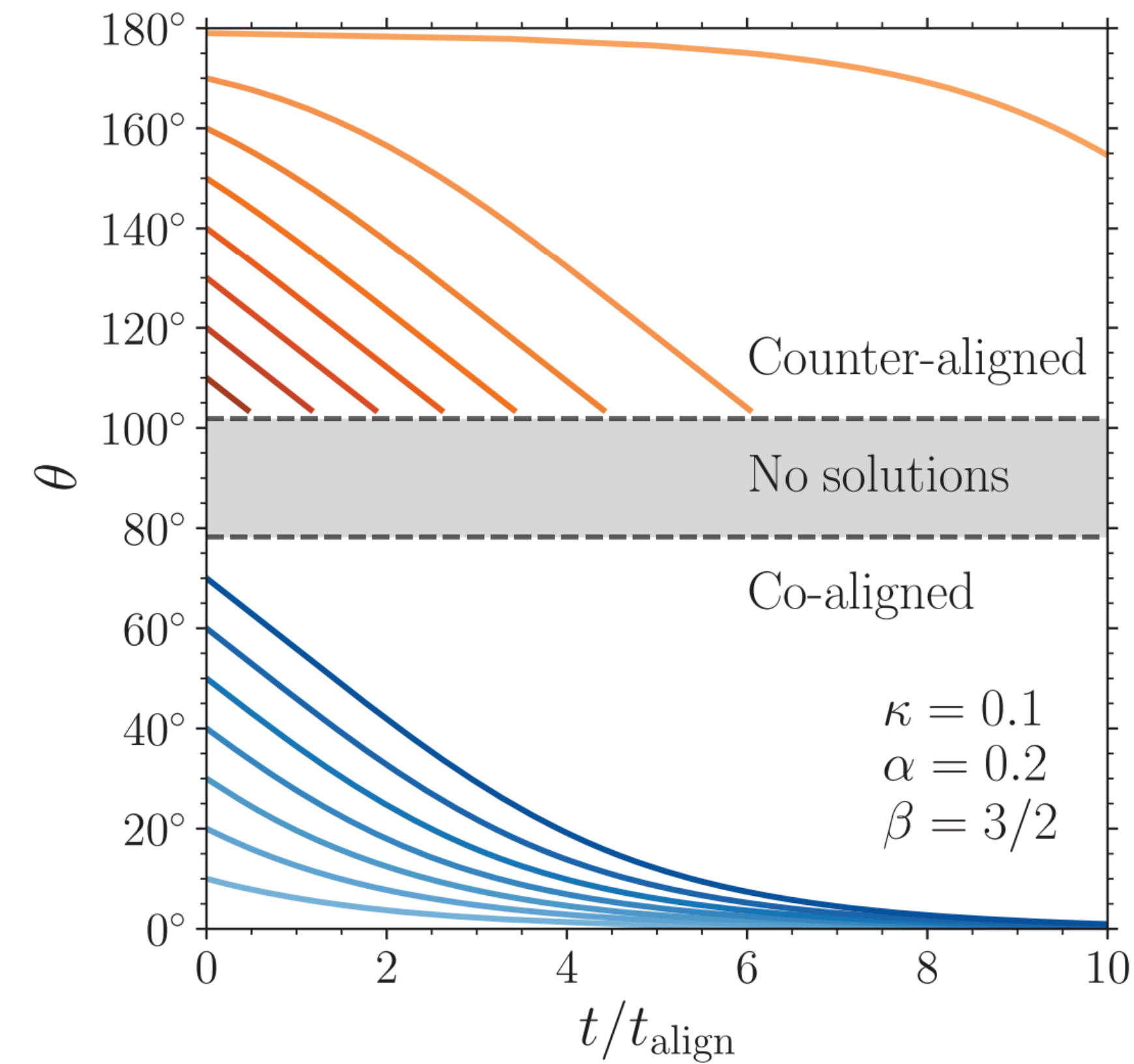
($\kappa=0$ reduces to a single black hole around the primary)

Solving the semi-analytic model

1. Choose disc parameters and parametrise the result with κ .
2. For a given angle, solve for the evolution of the surface density and angular momentum of the disc.
3. Repeat this for a range of inclinations.



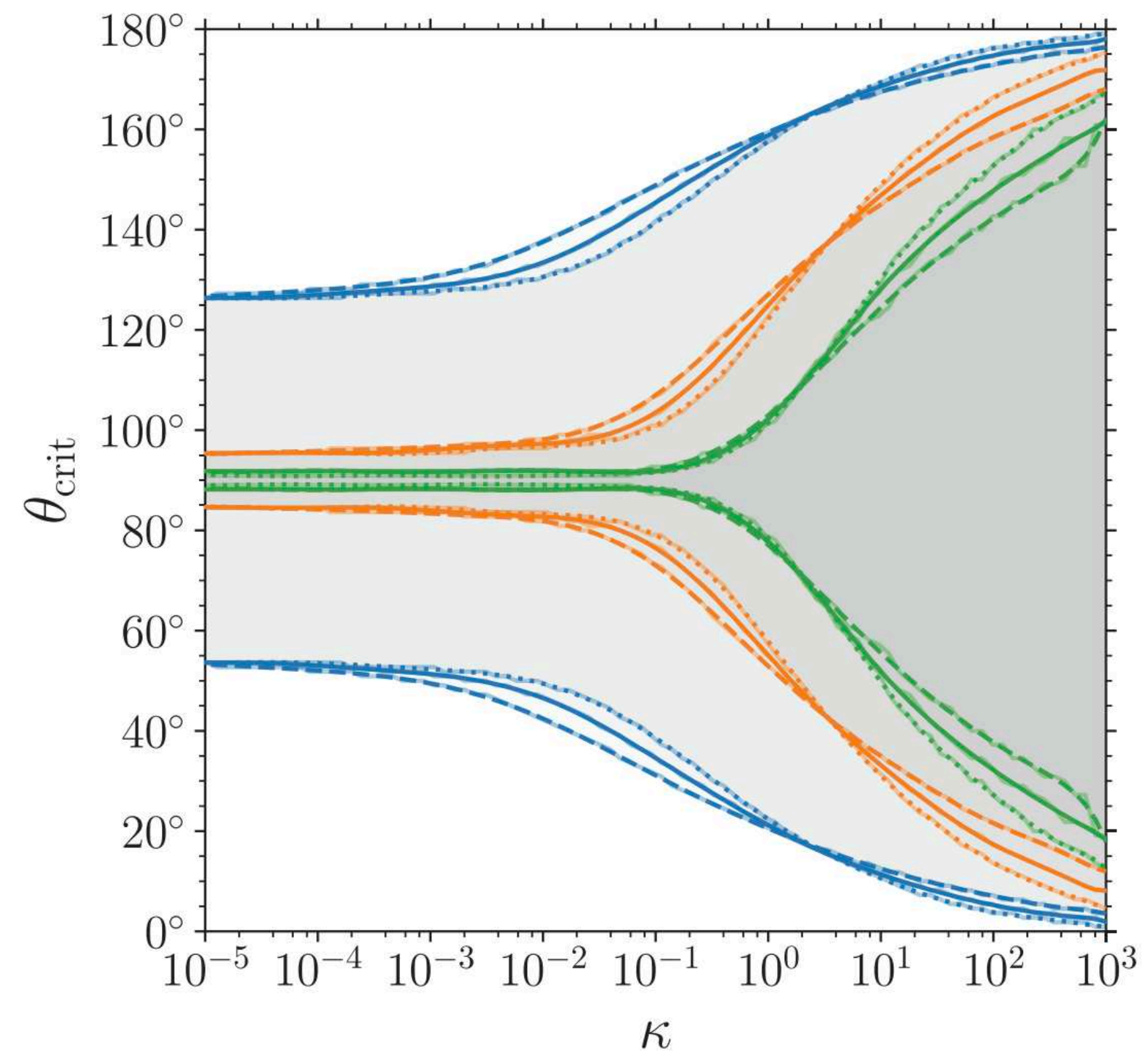
The break-down of the semi-analytic model



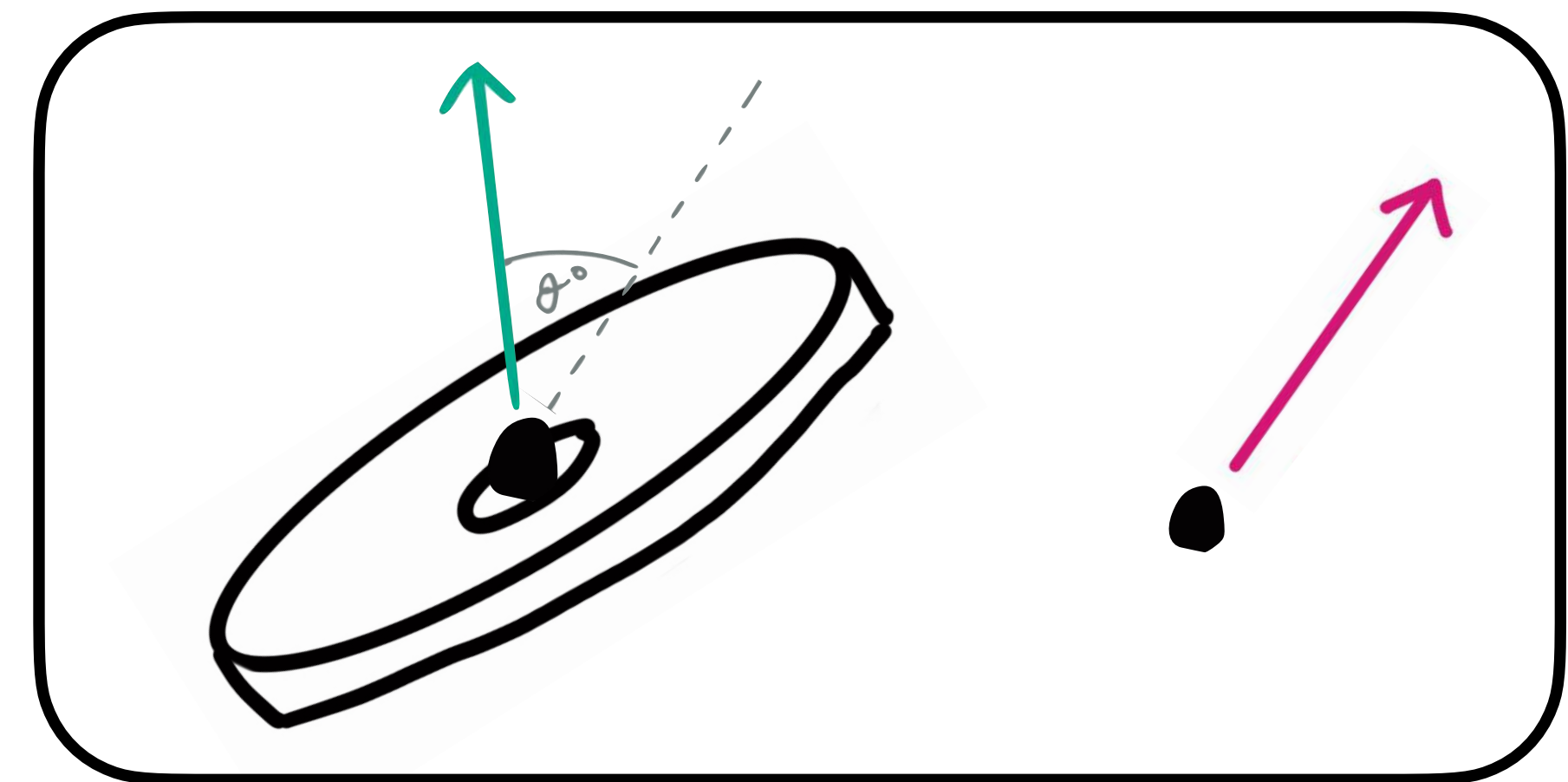
Our simulation suite

143 simulations in total:

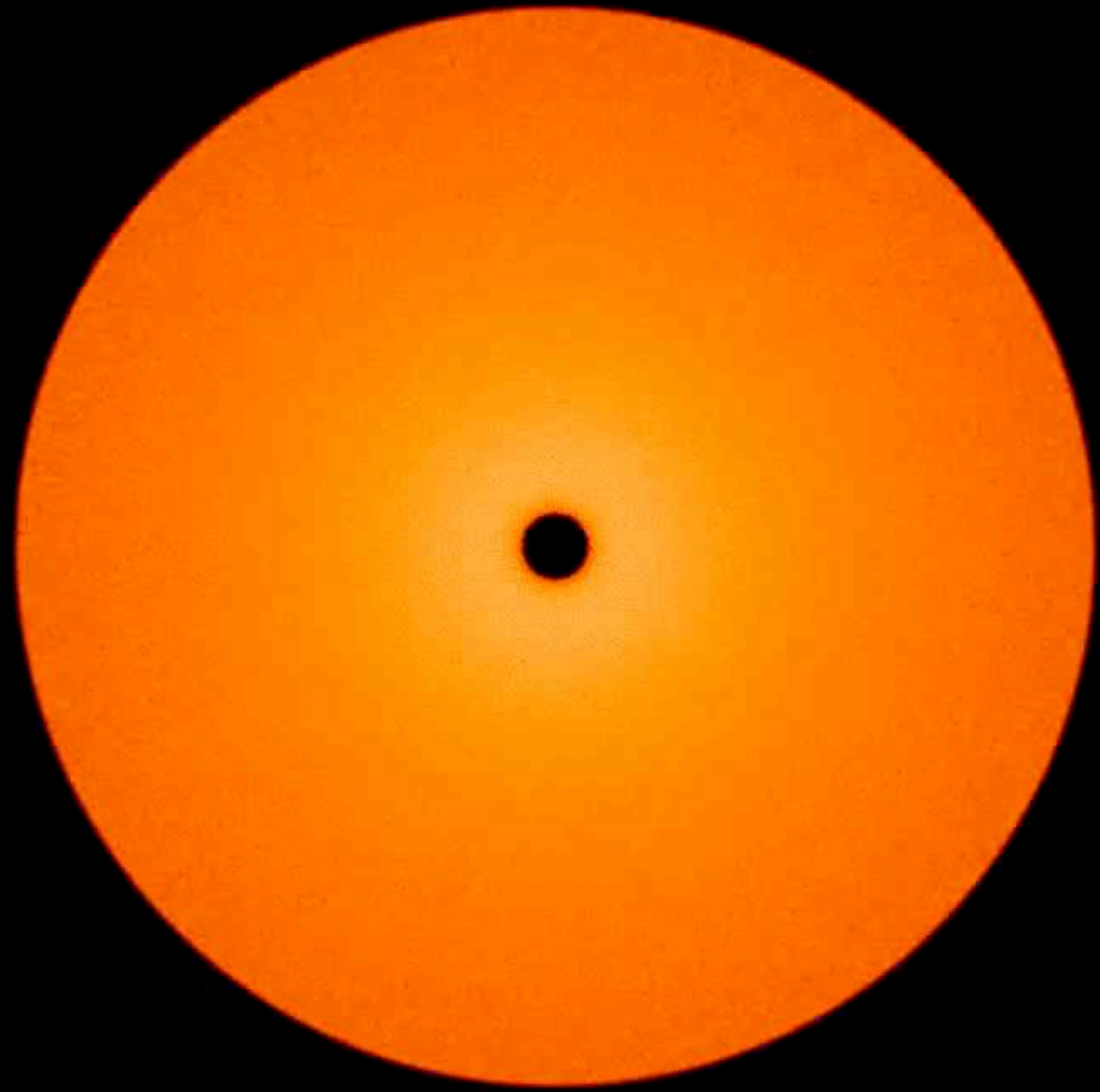
- 2 binary separations
 - Four viscosities
 - Four aspect ratios
 - Inclinations between 20 and 160 degrees
 - 12 κ values
-
- Use a fixed post-Newtonian potential
 - Secondary is much lower mass than primary
 - Black hole spin of $a=0.9$
 - Disc mass is 10^{-6} of the primary black hole
 - Minimum time of 150 orbits at R_{ref} (or 3 orbits of the outer binary)



Cover this
with these



0 orbits



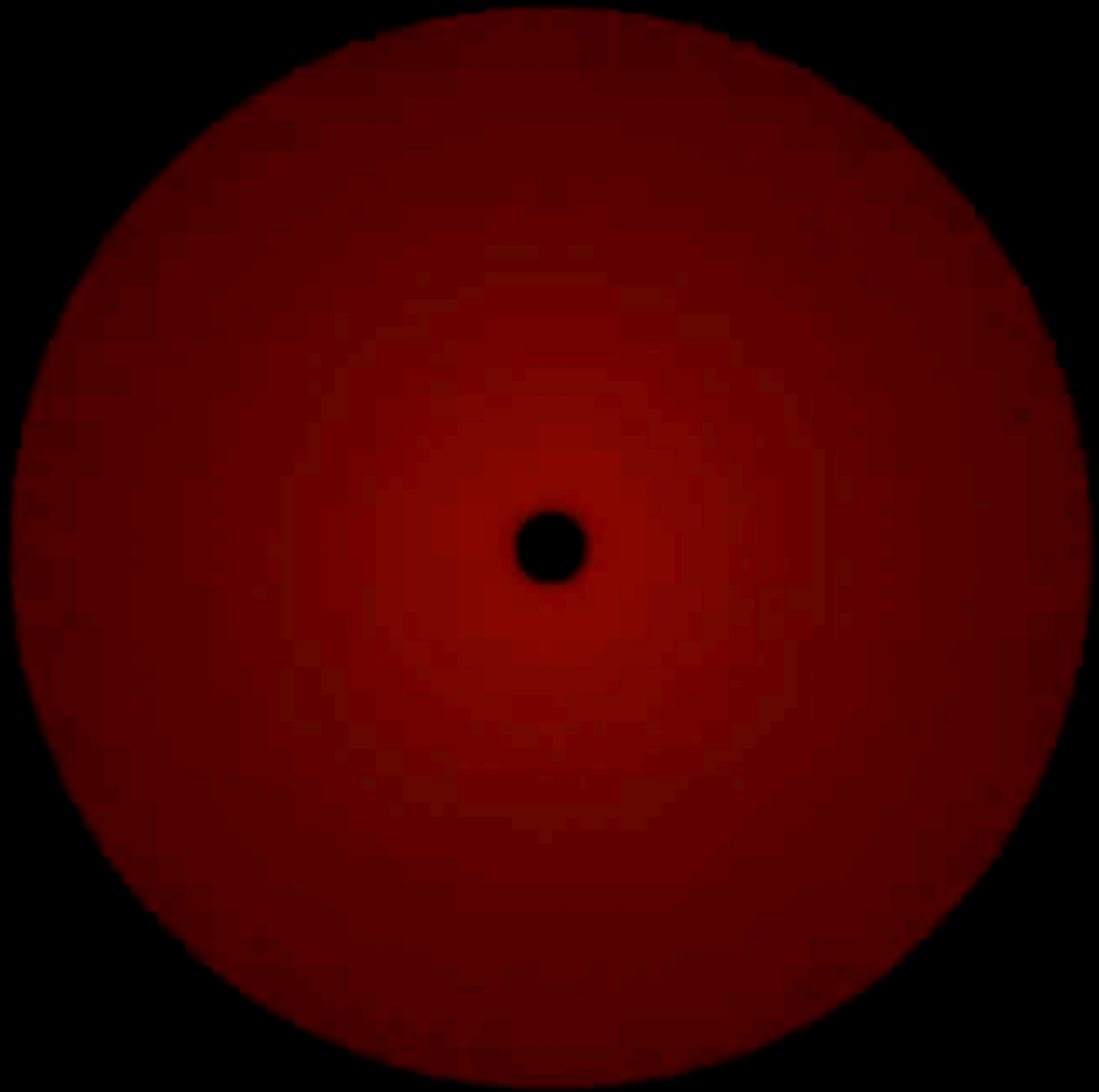
Warping

20 deg



Nealon, Ragusa, Gerosa, Rosotti & Barbieri 2022

0 orbits



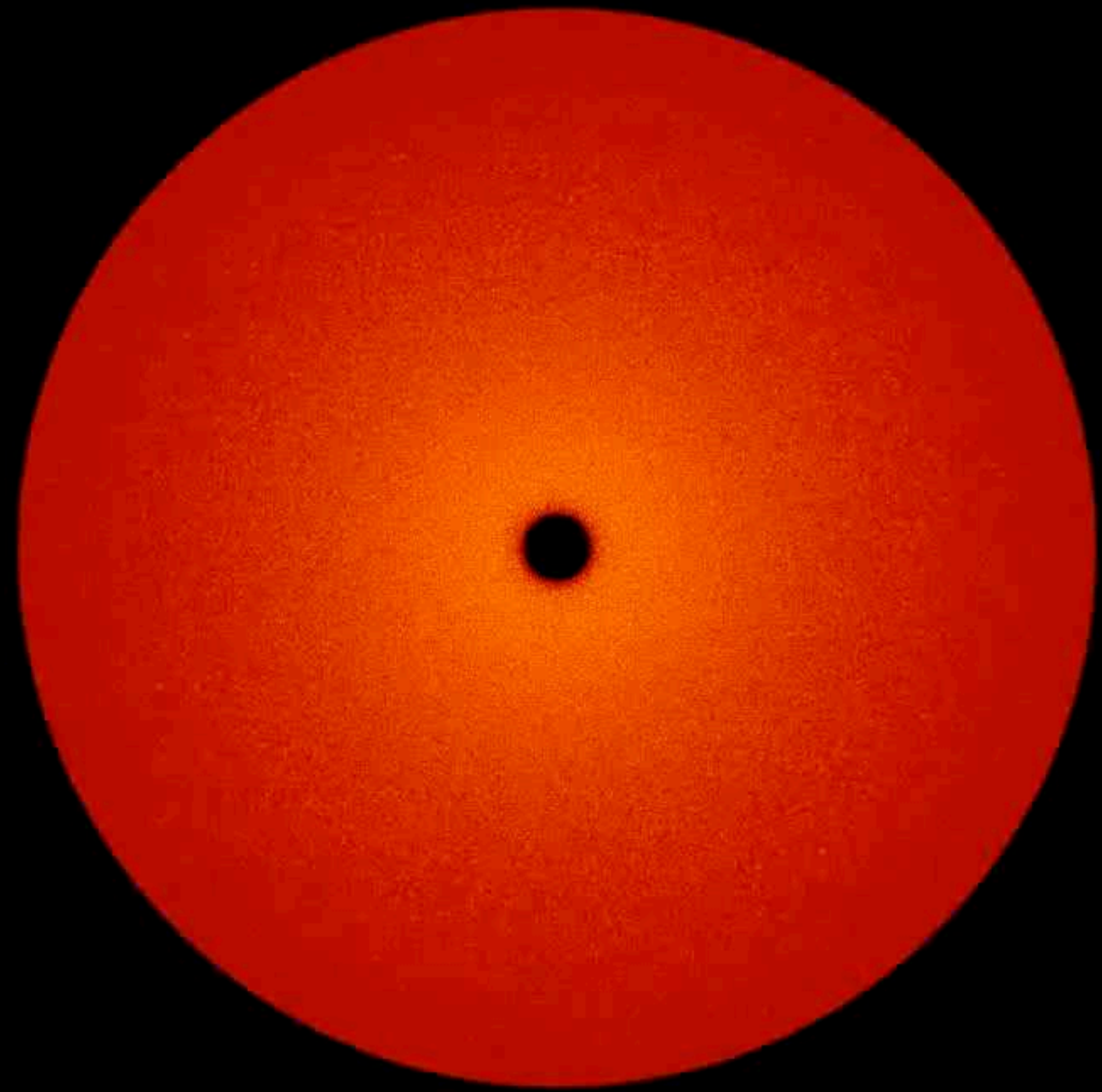
Breaking/tearing

60 deg

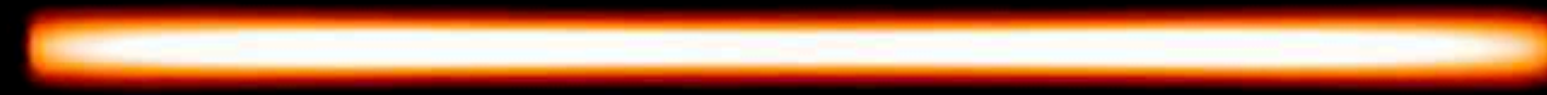


Nealon, Ragusa, Gerosa, Rosotti & Barbieri 2022

0 orbits

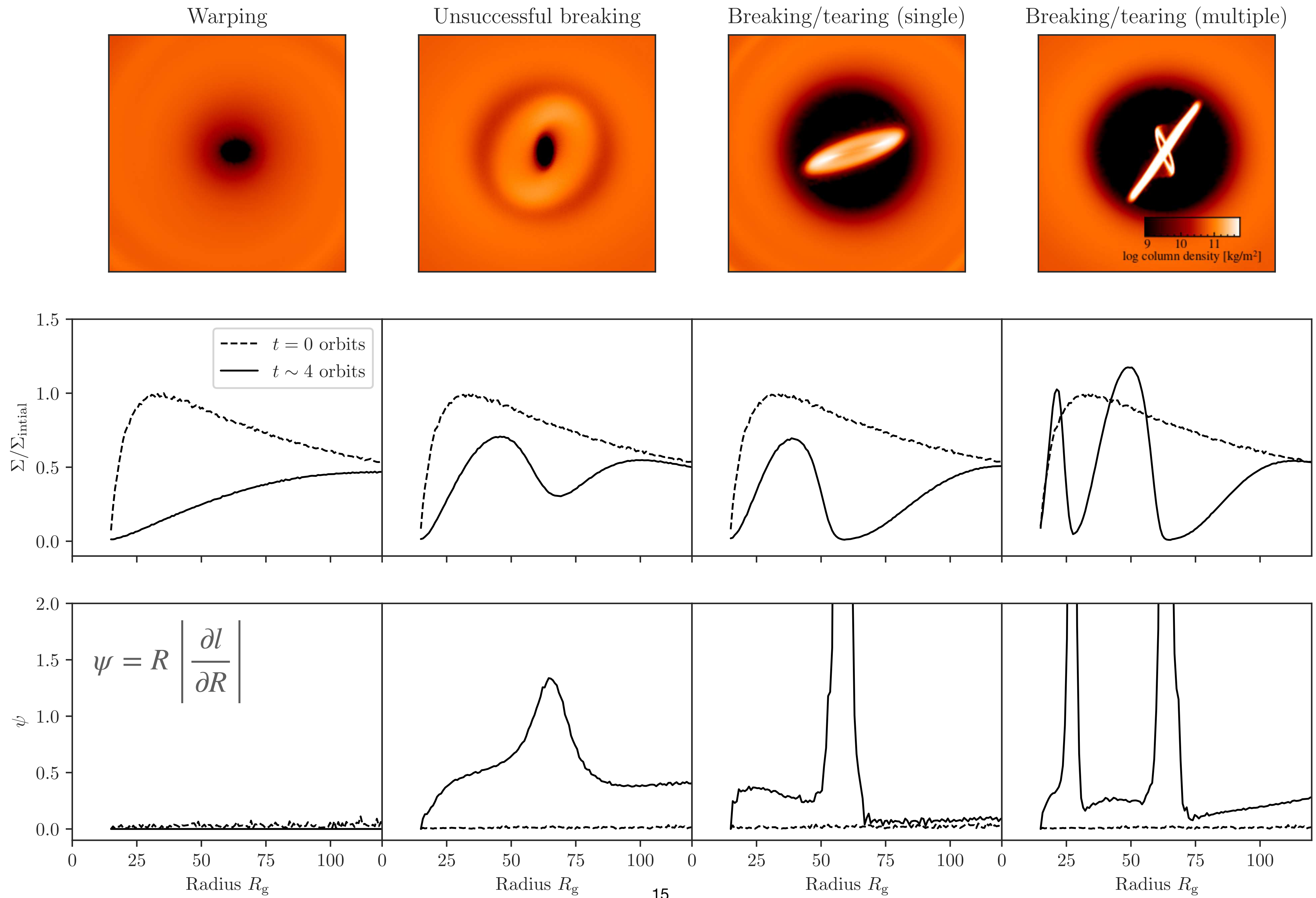


140 deg

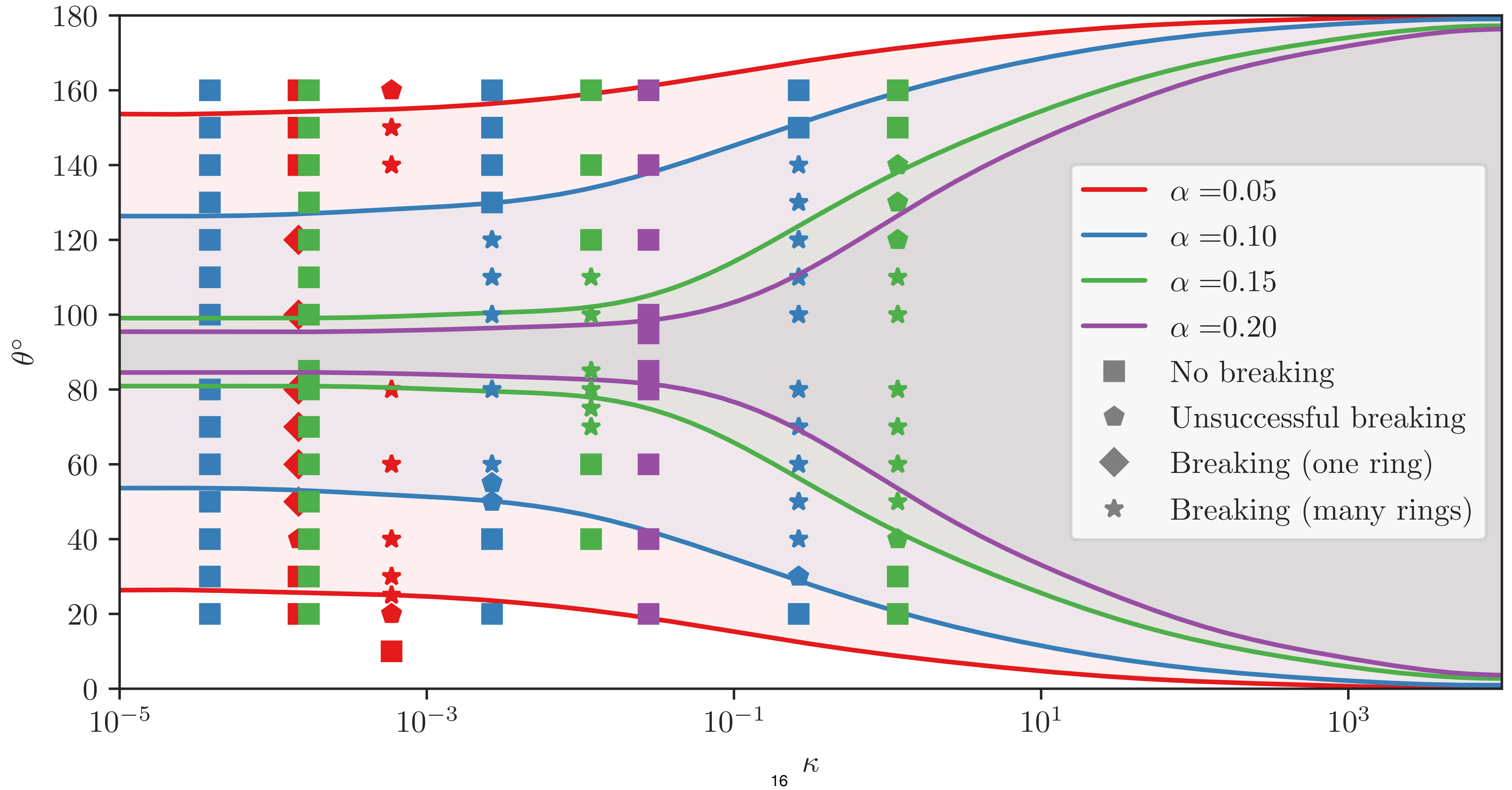


Unsuccessful breaking

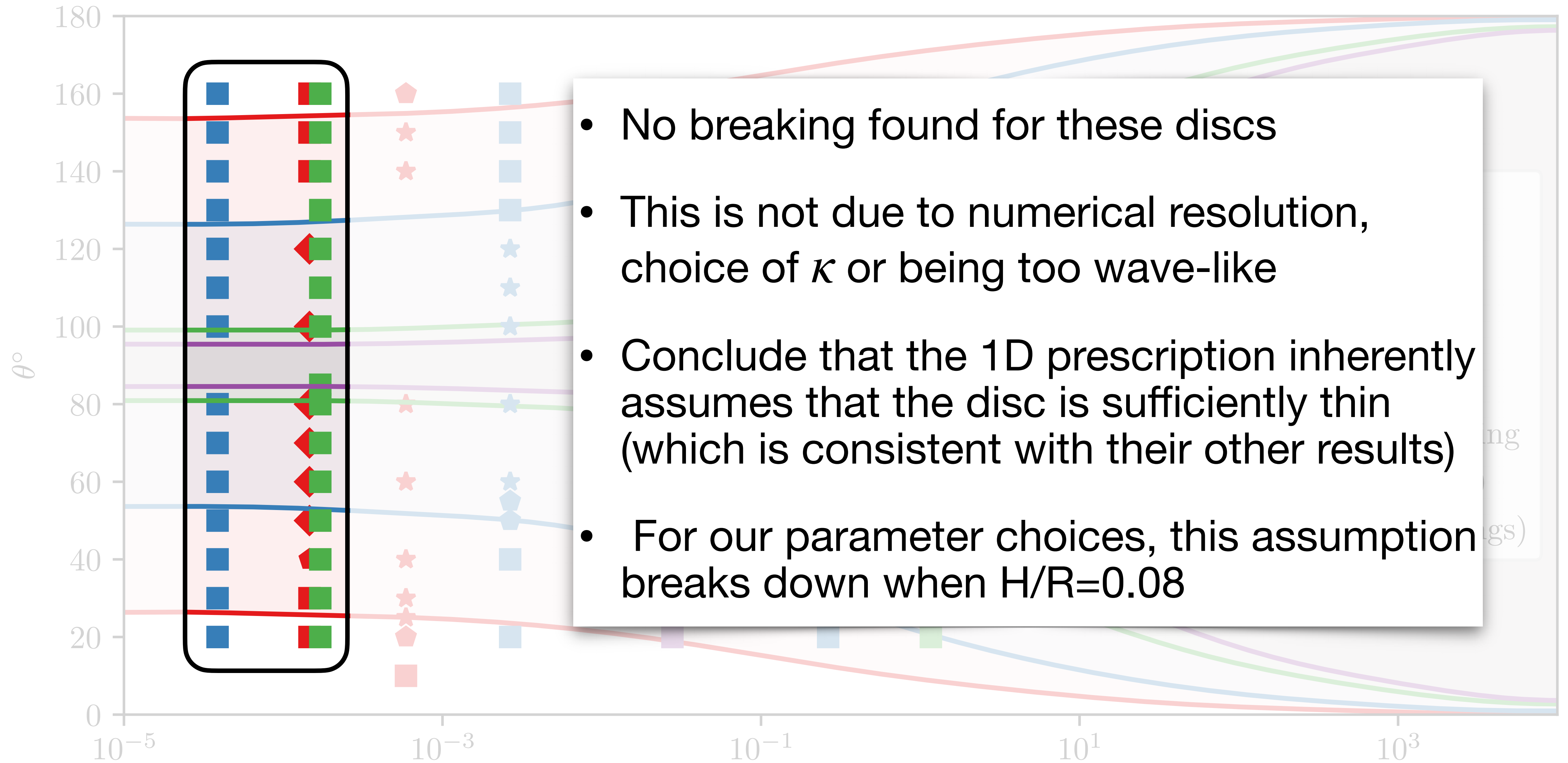
Defining a break



Comparing simulations and semi-analytic theory

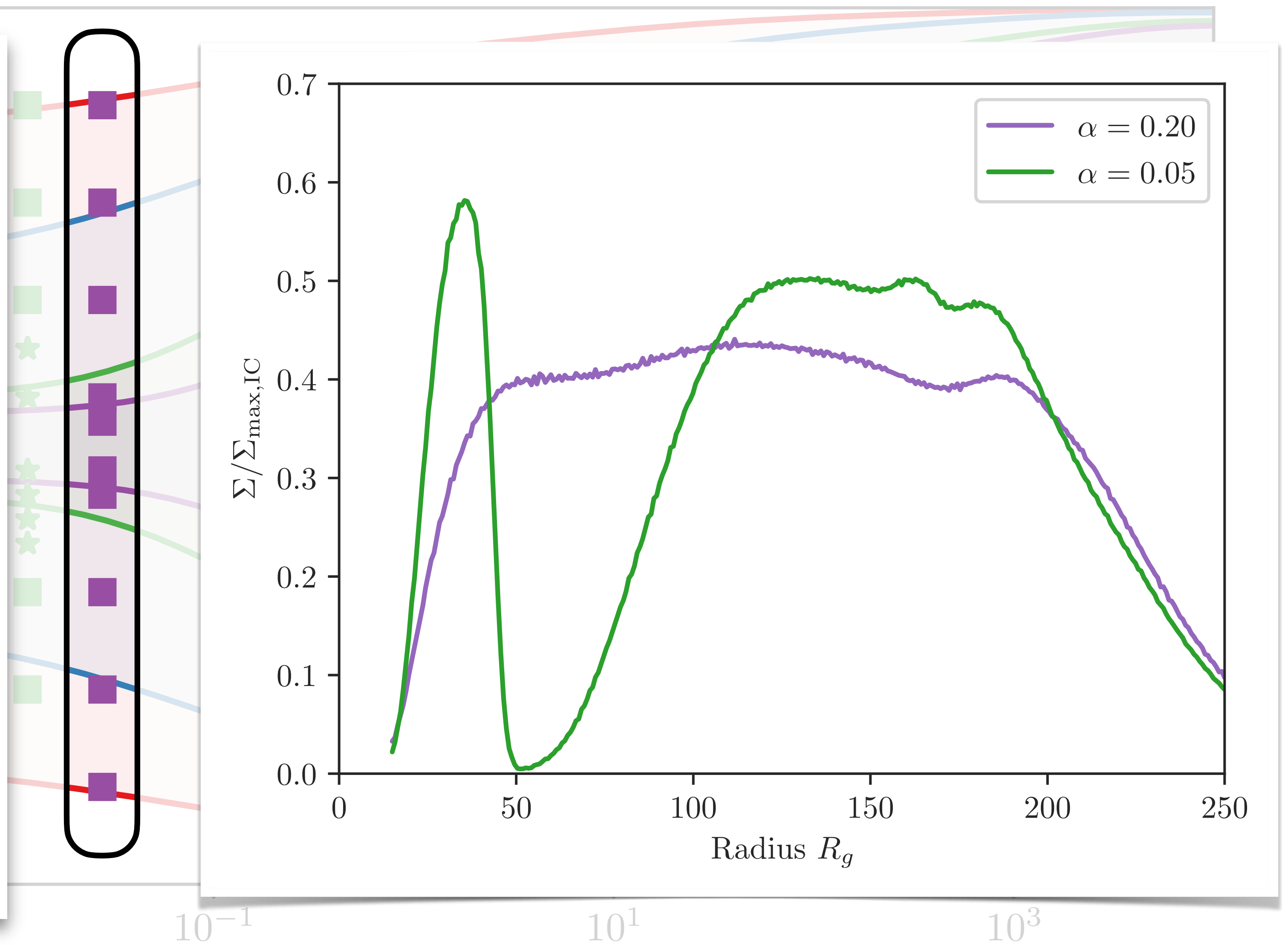


Thick discs don't break



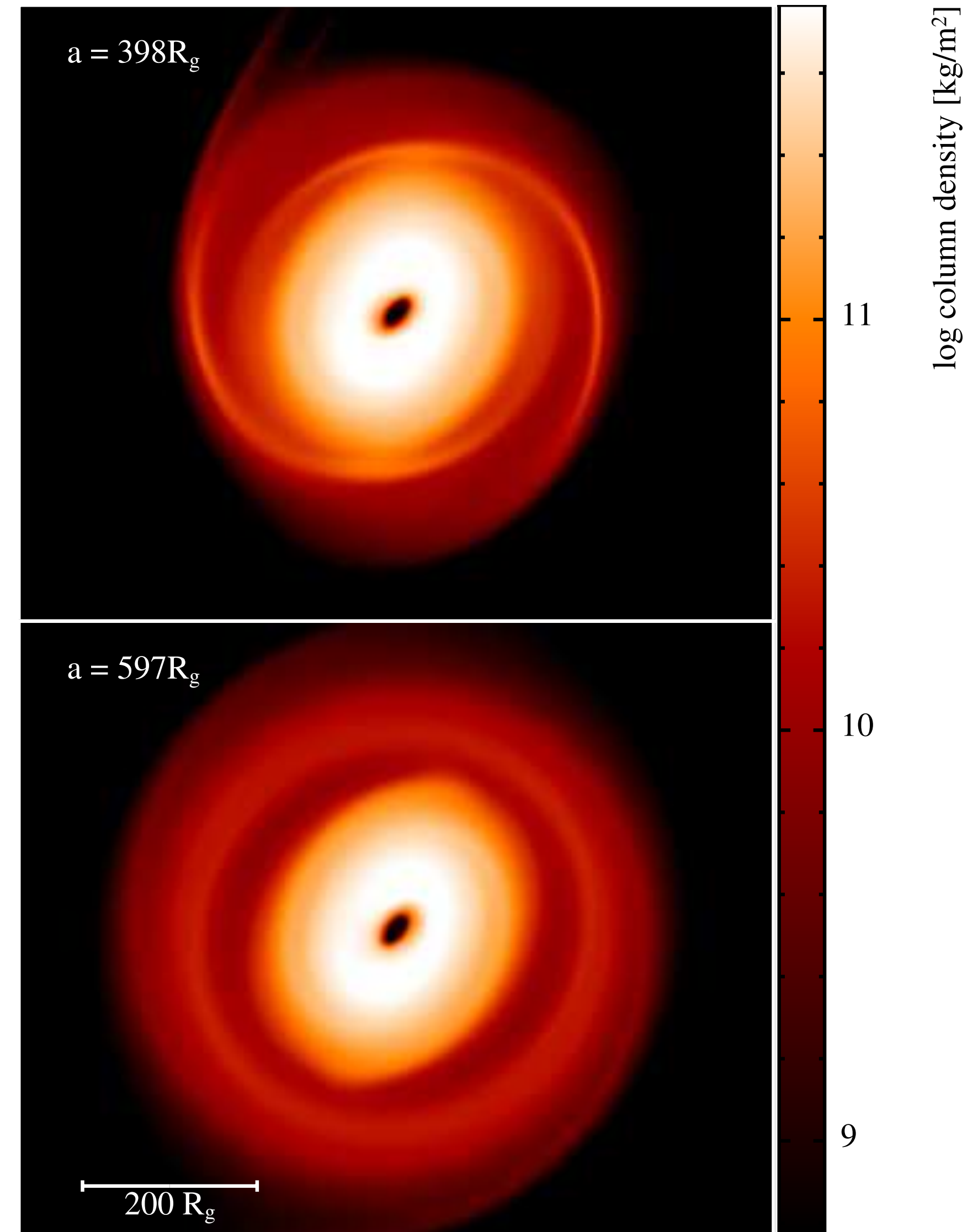
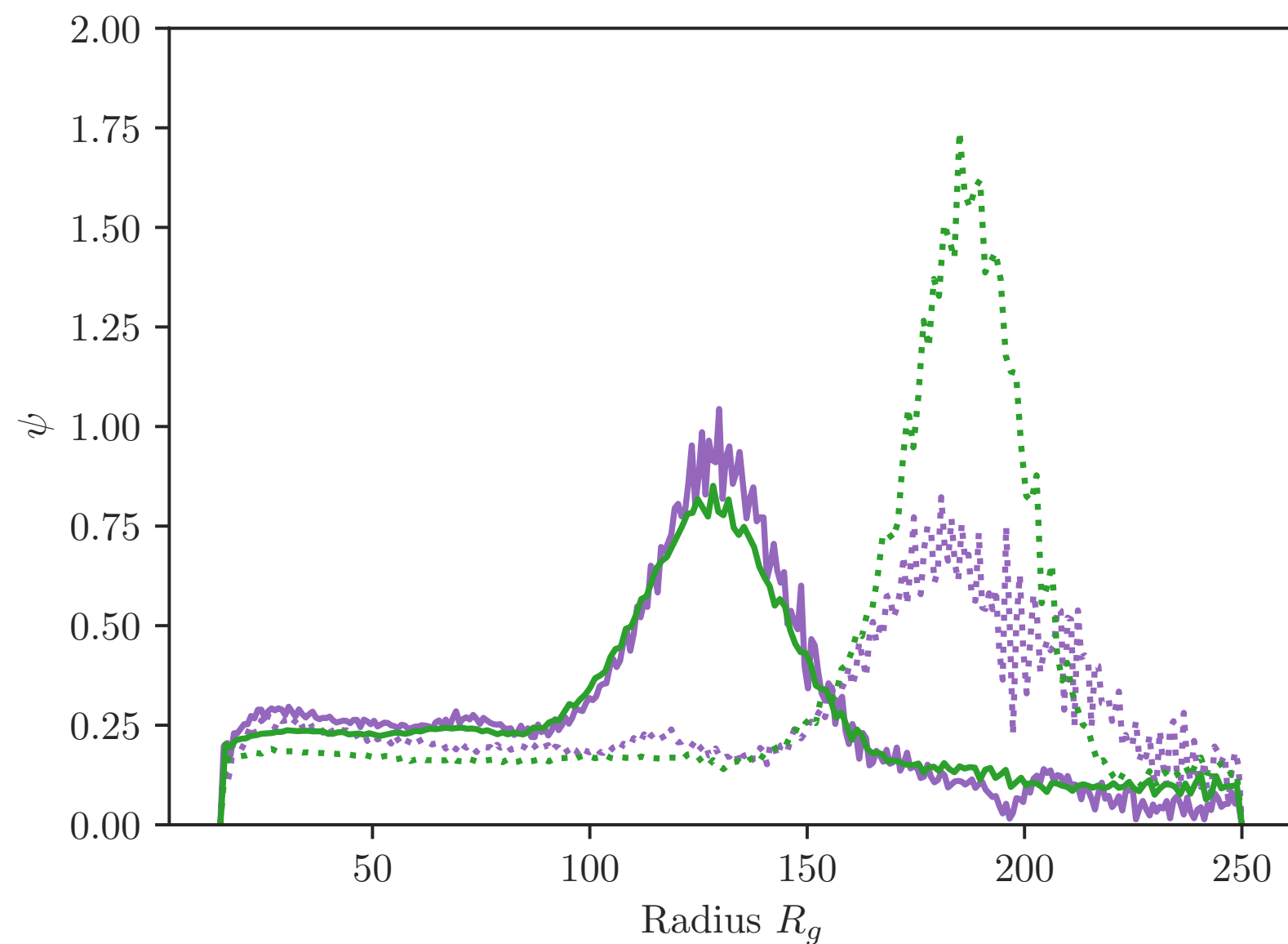
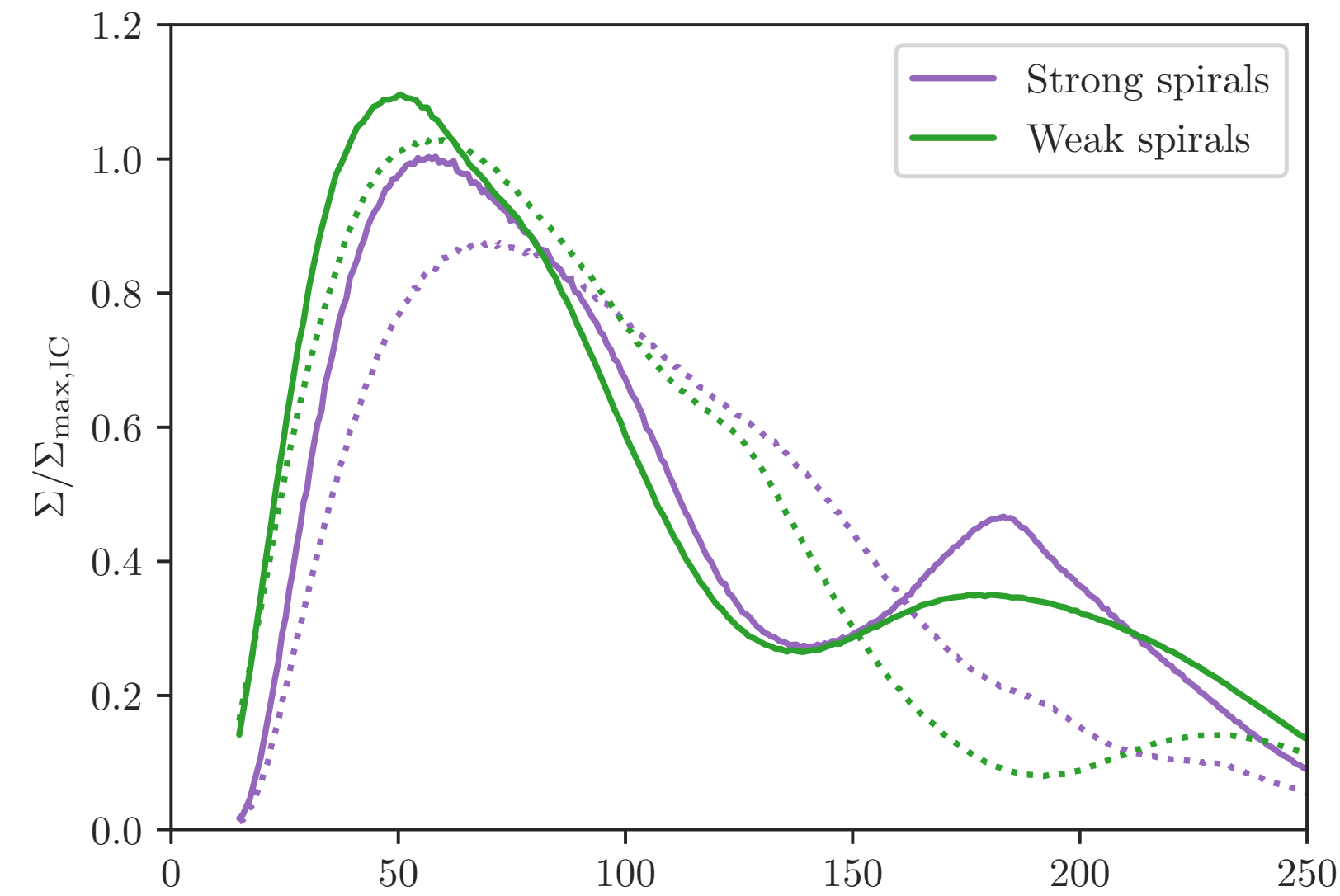
High viscosity discs don't break

- Again, not likely due to limitations of the simulations
- But looking at the surface density profile, we find that the inner disc accretes rapidly
- If breaking were to occur, this is where the inner ring would form
- Not enough material to make that inner ring -> no breaking



Spiral arms can prevent disc breaking

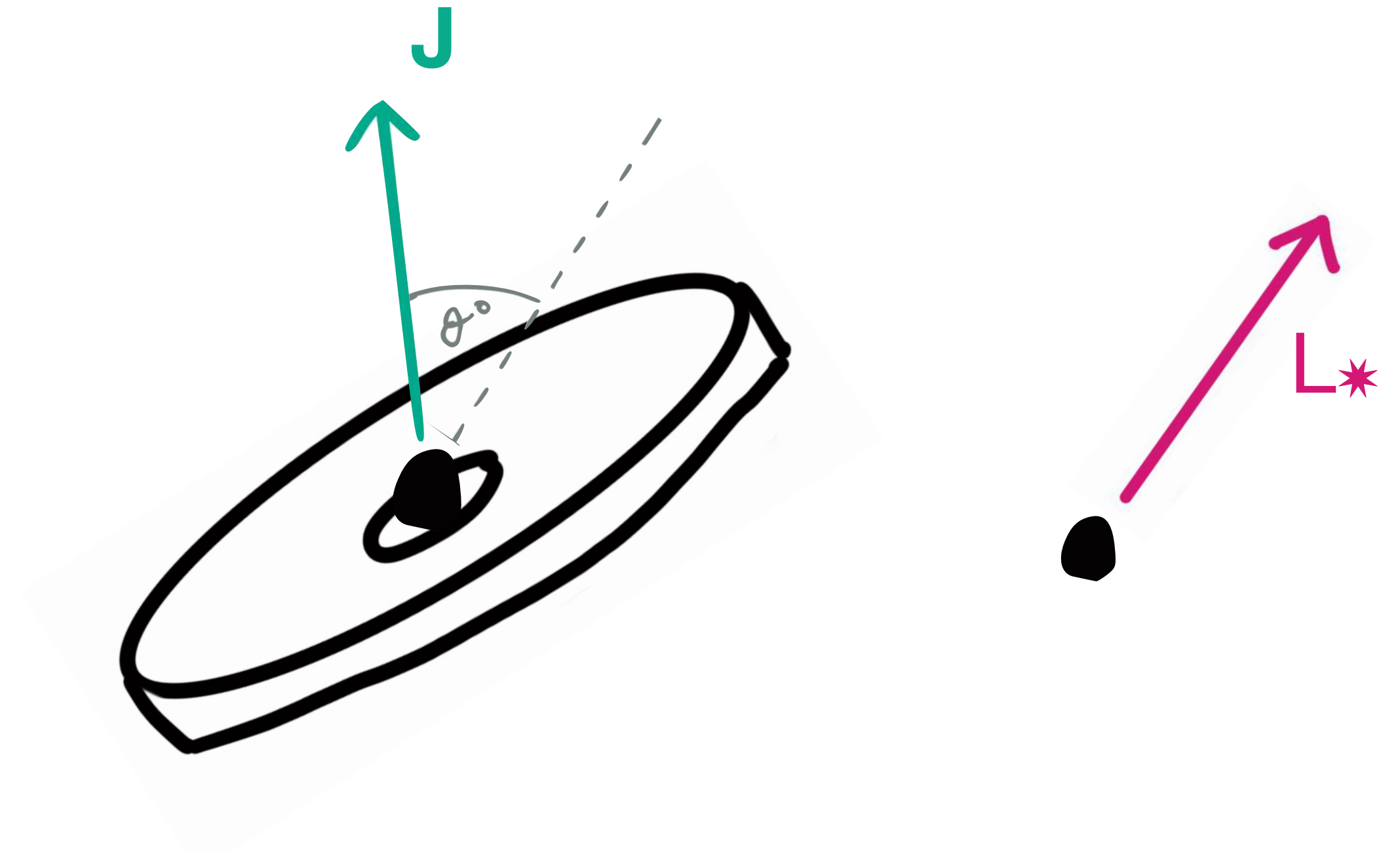
- We identify this one as unsuccessful breaking
- As the warp propagates out, it increases in amplitude until it hits the spiral arms
- Spiral arms increase the local disc viscosity, which in turn makes it harder to break the disc



Disc driven migration

$$\frac{d \cos \theta}{dt} = \frac{d\hat{\mathbf{J}}}{dt} \cdot \hat{\mathbf{L}}_*$$

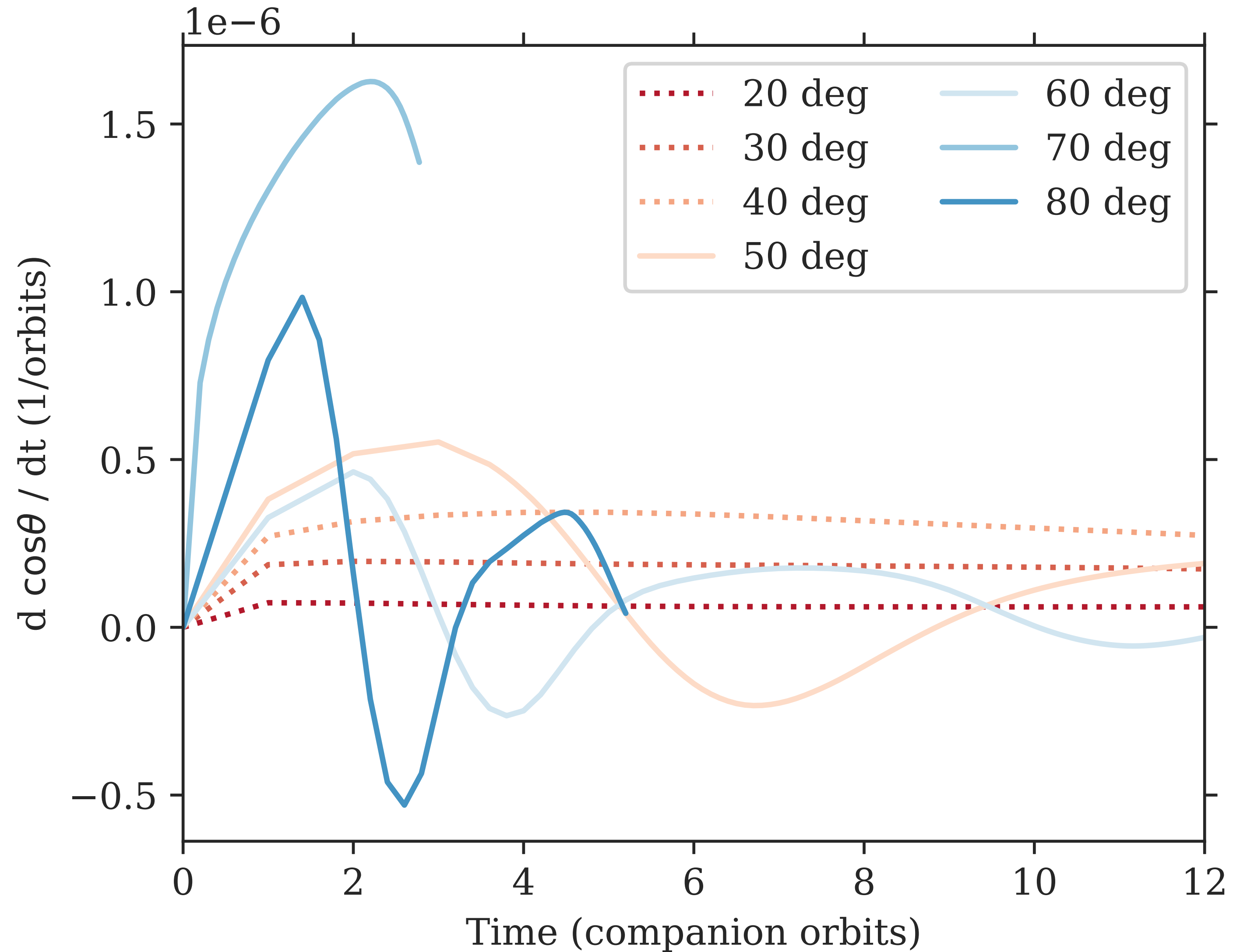
$$\frac{d\mathbf{J}}{dt} = - \int_{R_{\min}}^{R_{\max}} \frac{2G}{c^2} \frac{\mathbf{J} \times \mathbf{L}}{R^3} 2\pi R dR$$



1. What shape does the disc take?
2. How does this affect the alignment of the black holes?

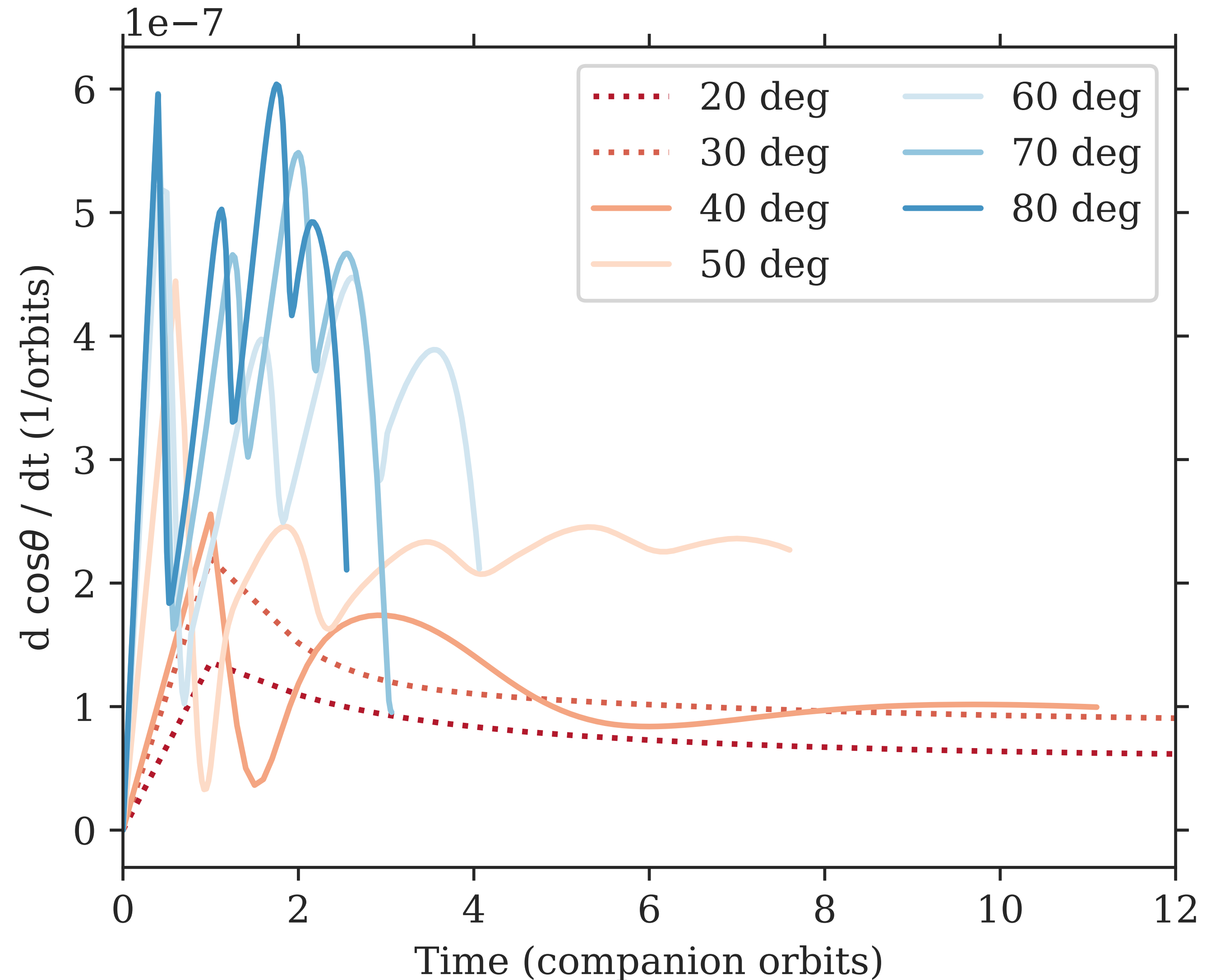
Disc-black hole misalignment

- Warping results in alignment
- Breaking with one ring results in oscillations, with occasional anti-alignment



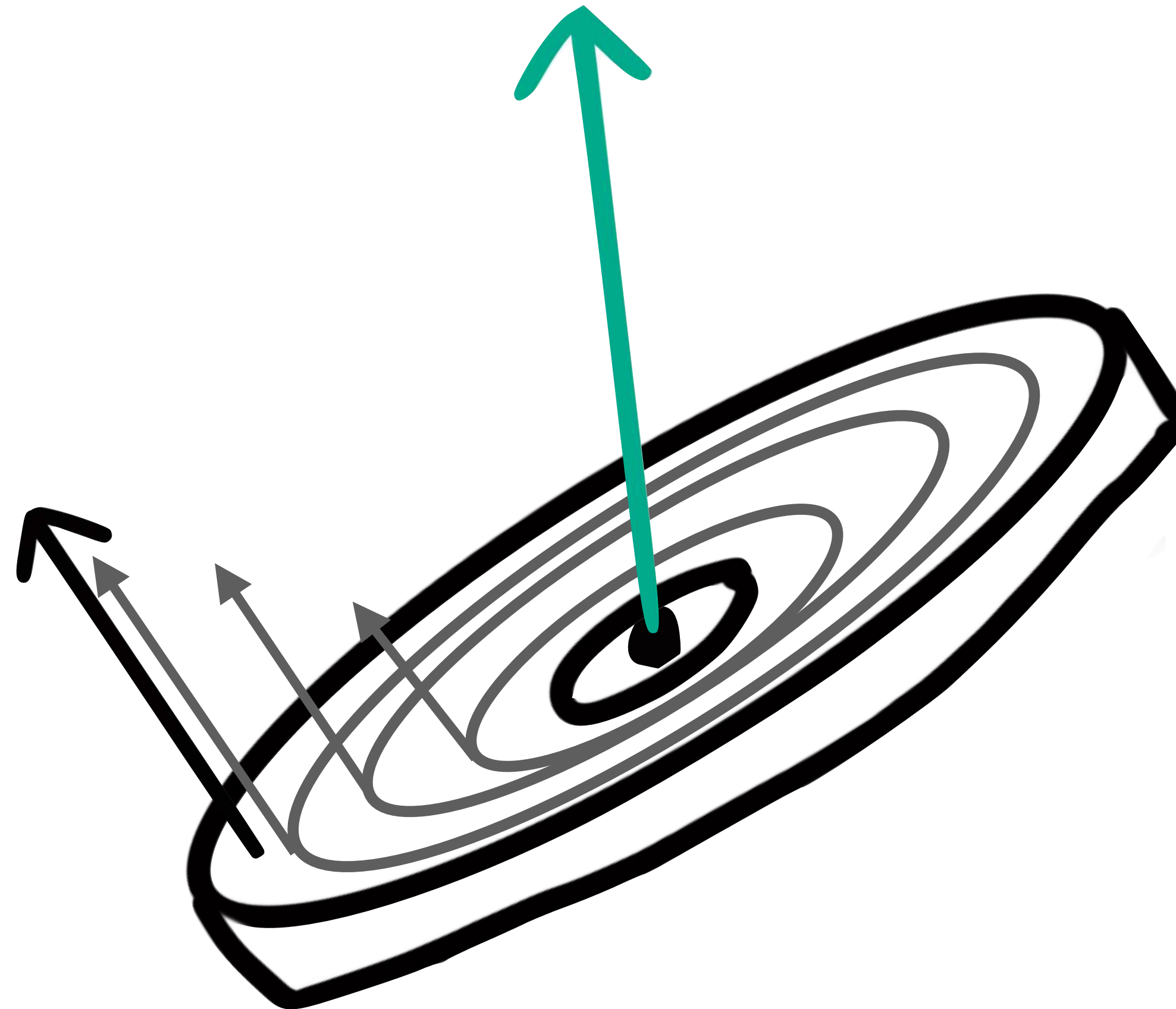
Disc-black hole misalignment

- Warping results in alignment
- Breaking with one ring results in oscillations, with occasional anti-alignment
- Breaking with multiple rings results in unpredictable oscillations



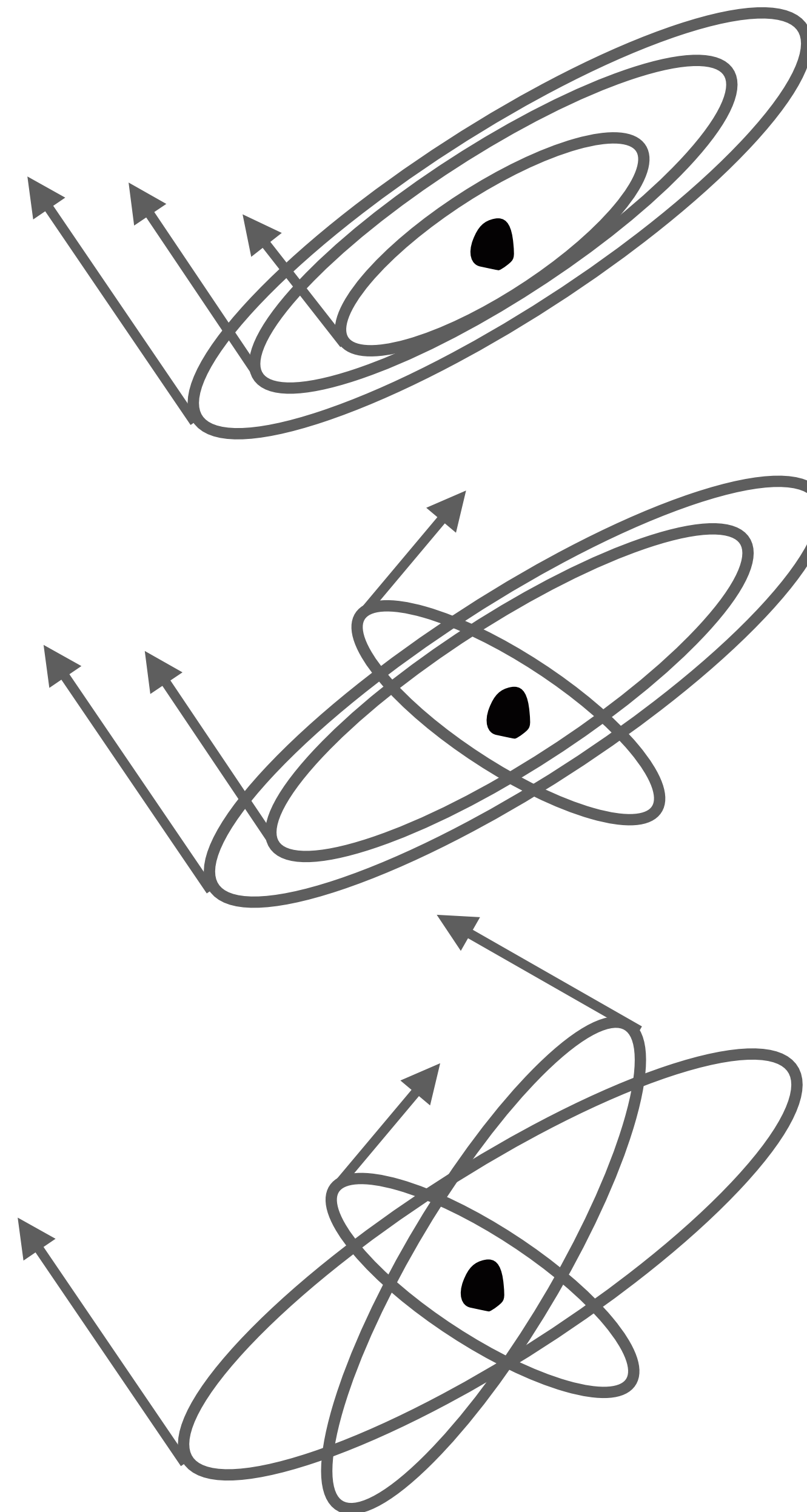
$$\frac{d\mathbf{J}}{dt} = - \int_{R_{\min}}^{R_{\max}} \frac{2G}{c^2} \frac{\mathbf{J} \times \mathbf{L}}{R^3} 2\pi R dR$$

- Warping results in alignment
- Breaking with one ring results in oscillations, with occasional anti-alignment
- Breaking with multiple rings results in unpredictable oscillations




$$\frac{d\mathbf{J}}{dt} = - \int_{R_{\min}}^{R_{\max}} \frac{2G}{c^2} \frac{\mathbf{J} \times \mathbf{L}}{R^3} 2\pi R dR$$

- Only warping results in alignment
- Breaking with one ring results in oscillations, with occasional anti-alignment
- Breaking with multiple rings results in unpredictable oscillations



Our results

- We find excellent agreement with the semi-analytic model of Gerosa et al. 2020
 - In the region where the semi-analytic model breaks down, we demonstrate that the discs are breaking
 - We have explained exceptions for very thick discs, large α and very large κ
 - Spiral arms can stabilise the disc against disc breaking
 - Disc breaking hinders (and in some cases prevents) alignment between the disc and the black hole
- 

- Followed up by Steinle & Gerosa et al. 2023, showing that there are likely to be distinct populations that should be observable