Radiation hydrodynamic simulations of common-envelope evolution

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### Heidelberg Old Bridge and Castle



## **Common-envelope evolution**



Donor Companion / accretor Dyna

A companion star enters the extended envelope of a giant star without co-rotation

Drag forces dissipate energy and angular momentum in the envelope

Dynamical plunge-in

Expelling the envelope leaves a much tighter binary orbit

Adapted from Schneider, Lau, Röpke (2025)



Envelope

ejection

Merger



Moe & Di Stefano 2017

### Stellar multiplicity is common

also, Offner+ 2022

Sana+ 2012

Binary interactions dominate massive star evolution



### Binary evolution tree

Han+ 2020





Ge 2020

### **Lau+** 2022a



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 $12~M_{\odot}$  red supergiant +  $3~M_{\odot}$  companion

Key questions:

- Can we simulate common-envelope ejection?
- What determines the post-commonenvelope orbital separation?

Current state:

- No 3D simulation has demonstrated complete envelope ejection selfconsistently on the dynamical timescale
- Mass ejection occurs on longer timescales
- Inclusion of recombination energy is key

Nandez+ 2015, Ivanova & Nandez 2016, Reichardt+ 2020, Sand+ 2020, **Lau**+ 2022a,b, Moreno+ 2022, González-Bolívar+ 2022



# Recombination energy is key

### **Lau+** 2022b



### **Lau+** 2022a



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### Need radiation transport!



# Radiation transport in common envelopes

### Ricker+ 2018 (IAU conference proceedings)



Flux-limited diffusion in the AMR code FLASH

Other implications of radiation transport:

- Transport away recombination energy
- Allows sensible lightcurves

Hatfull & Ivanova 2025



Flux-limited "emissiondiffusion" method in the SPH code StarSmasher

Lightcurve model for V1309 Sco

• Required for constructing a "realistic" giant star with envelope convection



# Radiation hydrodynamics scheme

Existing implementations of radiation transport in Phantom

- Two-temperature flux-limited diffusion following Whitehouse & Bate (2004) with an explicit solver
- $\rightarrow$  Tiny time steps in optically thin regions

Biriukov 2020 (PhD thesis)

$$\Delta t_{\rm rad} = C \frac{h^2 \rho \kappa}{c \lambda} \to C_{\rm Cour} \frac{h}{c} \ll t_{\rm dyn}$$

- Monte Carlo radiative transfer with MCFOST, CMacIonize
- $\rightarrow$  Too expensive in optically thick regions

Pinte+ 2006, 2009

- Port implicit radiative diffusion solver from SPHNG into Phantom, based on Whitehouse & Bate (2004) and Whitehouse+ (2005)
- Use Gauss-Seidel and backwards Euler method to solve gas and radiation energy equations
- Do not cap time step at  $\Delta t_{\rm rad}$
- Set accuracy tolerance for  $u, \xi$  solutions of  $10^{-6}$

- Optically thick regions converge within a few iterations
- Optically thin regions require considerably more (we set a maximum of 250)



# Radiation hydrodynamics scheme

Two-temperature flux-limited diffusion assuming LTE:

$$\begin{aligned} \frac{D\rho}{Dt} &= -\rho\nabla\cdot\mathbf{v} \\ \rho\frac{D\mathbf{v}}{Dt} &= -\nabla\rho + \frac{\kappa\rho}{c}\mathbf{F}_{\rm rad} + \rho\mathbf{\Pi}_{\rm shock} - \rho\nabla\Phi \\ \rho\frac{D\xi}{Dt} &= -\nabla\cdot\mathbf{F}_{\rm rad} - \nabla\mathbf{v}:\mathbf{P}_{\rm rad} + a_{\rm rad}c\kappa\rho(T_{\rm gas}^4 - T_{\rm rad}^4) \\ \rho\frac{Du}{Dt} &= -p\nabla\cdot\mathbf{v} + \rho\Lambda_{\rm shock} - a_{\rm rad}c\kappa\rho(T_{\rm gas}^4 - T_{\rm rad}^4) \\ \end{aligned}$$

$$\begin{aligned} \text{Mihalas \& Mihalas 1984, Turner \& Stone} \end{aligned}$$

Whitehouse & Bate 2004

Ideal gas EoS 
$$p = \frac{\rho \Re T_{\text{gas}}}{\mu}$$
  $u = \frac{p}{(\gamma - 1)\rho}, \ \gamma = 5/3$   
Radiation energy  $\xi = a_{\text{rad}} T_{\text{rad}}^4 / \rho$ 

Radiative flux  $\mathbf{F} = - \stackrel{e_{H}}{\longrightarrow} \nabla E \longleftarrow$  Radiation energy density, κρ  $E = \rho \xi$ 

 $\lambda \rightarrow 1/3$ Optically thick limit Flux limiter

> $|\mathbf{F}| \rightarrow cE$ Optically thin limit

> > Levermore & Pomraning 1981

Opacity tables from the MESA stellar evolution code implemented in *Phantom* by Reichardt+ 2020

Tables combine OPAL opacities for the high-temperature regime with SCVH opacities for the low-temperature regime

> Rogers+1996, Rogers & Nayfonov 2002, Saumon+ 1995

2001,

Features/assumptions:

- $\xi$  diffuses along  $-\nabla T_{rad}$
- $T_{\rm gas}$  and  $T_{\rm rad}$  are decoupled in the optically thin limit
- LTE allows source term to be written as Planck function (but not thermal equilibrium)
- Gray opacity
- No ionisation/recombination energy, but includes opacity changes









# Test problems

Gas heating/cooling Biriukov 2020 (PhD thesis)  $10^{10}$ . Gas cooling  $10^{9}$  $10^{8}$  $10^{7}$ erg g  $10^{6}$ n $10^{5}$ 2. Gas heating  $10^{4}$  $10^{3}$  $10^{2}$  $10^{-17}$  $10^{-15}$  $10^{-13}$  $10^{-11}$  $10^{-9}$  $10^{-7}$ t / s

Solid lines: Explicit solver

Markers: Implicit solver (40 time steps)

Diffusion of sinusoid

Biriukov 2020 (PhD thesis)



Solid lines: Analytical solution at different times

Markers: Implicit solver



- Can be reproduced using new option in SETUP=shock
- $\xi_{\text{left}} = 10\xi_{\text{right}}$







# Setup

- Initial conditions are identical to past adiabatic simulations  $\bullet$ (**Lau+** 2022a,b)
- Donor star is based on a MESA red supergiant model:  $\bullet$ 
  - $M = 12 M_{\odot}, R = 618 R_{\odot}, L = 1.68 \times 10^{38} \text{ erg s}^{-1}$
  - Dense helium core is replaced with a point mass that only interacts via cubic-spline-softened gravity
  - Envelope resolved with 2 million SPH particles
- Represent companion using a  $3~M_{\odot}$  point mass in a circular, Keplerian orbit

We do not drive envelope convection in the pre-commonenvelope donor, and prescribe a flat-entropy initial profile that is convectively stable

**Lau**+ 2022a



$$\frac{L\Delta t_{\text{inspiral}}}{E_{\text{bind}}} = 0.022 \left(\frac{L}{1.7 \times 10^{38} \text{ erg s}^{-1}}\right) \left(\frac{\Delta t_{\text{inspiral}}}{5 \text{ yr}}\right) \left(\frac{E_{\text{bind}}}{1.2 \times 10^{48} \text{ erg}}\right)$$

### Face-on density slice



### Edge-on density slice



**Lau+** 2025



 No significant differences in the post-plunge-in orbital separation



Solid lines: Unbound  $\iff e_k + e_p > 0$ Dashed lines: Unbound  $\iff e_k + e_p + e_{th} > 0$ 

Radiative diffusion significantly suppresses envelope ejection



**Lau+** 2025



No significant differences in the post-plunge-in orbital separation



Dashed lines: Unbound  $\iff e_k + e_p + e_{th} > 0$ 

Radiative diffusion significantly suppresses envelope ejection



### Edge-on <u>density slice</u>

### Adiabatic

Radiative diffusion

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### Clear bipolar morphology driven by fast outflows near the central binary

Strong suppression of bipolar morphology due to fall back of ejecta. Obtain weak, intermittent plumes instead







### Face-on view



### Define local diffusion time,

$$t_{\text{diff}} := \frac{\kappa \rho}{c \lambda} \left( \frac{E}{|\nabla E|} \right)^2,$$

as the diffusion time across the temperature scale height

- Steep drop in  $t_{\rm diff}$  from ~ 10 to  $\lesssim 10^{-4}$  yr over poorly resolved region
- Steep drop in t<sub>diff</sub> is associated with opacity drop from H recombination
- Bound material has very short  $t_{\rm diff}$



# Wrapping up: Room for improvement

<u>Recombination energy</u>

- Add recombination energy, making substantial progress towards settling the debate on its relevance Sabach+ 2017, Grichener+ 2018, Ivanova 2018, Soker+18
- Cannot use MESA EoS tables directly, because radiation energy must be separated out from the total internal energy
- Current progress: Use Ryo Hirai's fits analytical treatment of ionisation physics (gas + radiation + recombination) EoS, ieos=20), including new fits of  $c_V$  accounting for rotational and vibrational degrees of freedom of H<sub>2</sub>

### <u>Convection and optically thin radiation transport</u>

- A realistic 3D giant star must have convection driven by photospheric cooling
- Photospheric cooling is not correctly captured due to the unresolved photosphere and an initial lack of SPH particles above the photosphere to radiate into

### Ma+ 2025: $10 \text{ M}_{\odot}$ red supergiant



Full 4π envelope convection with resolved photosphere in AREPO using Implicit Discrete **Ordinates Radiation** Transport







# Astrophysical implications

### <u>Towards non-adiabatic envelope ejection</u>

• A substantial fraction of the envelope is removed on timescales much longer than the dynamical timescale, a so-called *self-regulated phase* 

Meyer & Meyer-Hofmeister 1979, Podsiadlowski 2001

 Post-common-envelope binaries with distant tertiary companion support ejection timescales of  $10^2 - 10^4$  yr

Michaely & Perets 2019, Igoshev+ 2020

Traditional "energy formalism" breaks down

Hirai & Mandel 2022

### Implications of circumbinary material

- Circumbinary disk torques excite binary eccentricity and tighten orbit
- Circumbinary material could provide circumstellar material for *interacting supernovae* (e.g., SN 2014c)

### <u>Circumbinary disk + bipolar outflows in MHD simulations</u> Vetter+ (inc. Lau) 2024, 2025



3D rendering of magnetic field lines in bipolar outflow

### Gagnier & Pejcha 2025



MHD simulation focusing on post-CE bound ejecta

### Ondratschek+ 2022



### Comparison of ejecta with the Calabash nebula



Parameter study of circumbinary disk torques





# Conclusions

- phase in the formation of many astrophysical objects
- We performed one of the first 3D radiation hydrodynamic
- plunge-in orbital separation
- envelope is ejected after the plunge-in

Common-envelope evolution is a crucial but poorly understood

simulations assuming flux-limited diffusion and grey opacities

 Radiative diffusion significantly suppresses bipolar outflows, reduces the amount of ejected mass, but does not alter the post-

• Our results support the emerging paradigm of a multi-step common-envelope evolution where a significant amount of the